

## Regular Perturbation Theory (Differential Eqns)

Seek solution expansions of the form

$$(1) \quad y(t, \varepsilon) = y_0(t) + \varepsilon y_1(t) + \varepsilon^2 y_2(t) + \dots$$

to problems like

$$(2) \quad F(y'', y', y, t, \varepsilon) = 0 \quad 0 < \varepsilon \ll 1$$

May also include initial conditions

$$(3) \quad y(0, \varepsilon) = a \quad y'(0, \varepsilon) = b$$

where  $a, b$  could also depend on  $\varepsilon$ .

For instance

$$y'' + \varepsilon \sin t y' + y = y^2$$
$$y(0) = 1 \quad y'(0) = 2 + \varepsilon$$

The method of solution involves substituting the expansion (1) into the differential equation and initial conditions, expanding and then collecting like powers of  $\varepsilon$ .

EXAMPLE

$$y'' + y = \varepsilon y^2 \quad y(0) = 1 \quad y'(0) = 0$$

Assume

$$y(t, \varepsilon) = y_0(t) + \varepsilon y_1(t) + \varepsilon^2 y_2(t) + \dots$$

Expand initial conditions

$$y(0, \varepsilon) = y_0(0) + \varepsilon y_1(0) + \dots = 1$$

$$y'(0, \varepsilon) = y_0'(0) + \varepsilon y_1'(0) + \dots = 0$$

From which  $y_k'(0) = 0$ ,  $y_0(0) = 1$ ,  $y_k(0) = 0$  for  $k \geq 1$ .

Expand differential equation

$$(y_0'' + \varepsilon y_1'' + \dots) + (y_0 + \varepsilon y_1 + \dots) = \varepsilon (y_0 + \varepsilon y_1 + \dots)^2$$

Collecting powers of  $\varepsilon$  and using initial conditions we arrive at a sequence of initial value problems

$$(1) \quad y_0'' + y_0 = 0 \quad y_0(0) = 1 \quad y_0'(0) = 0$$

$$(2) \quad y_1'' + y_1 = y_0^2 \quad y_1(0) = 0 \quad y_1'(0) = 0$$

Clearly  $y_0(t) = \cos t$ . Then (2) becomes

$$y_1'' + y_1 = \cos^2 t = \frac{1}{2} + \frac{1}{2} \cos 2t$$

Using undetermined coefficients (method)

$$y_1(t) = \frac{1}{2} - \frac{1}{3} \cos t - \frac{1}{6} \cos 2t$$

So that

$$y(t, \varepsilon) = \cos t + \varepsilon \left( \frac{1}{2} - \frac{1}{3} \cos t - \frac{1}{6} \cos 2t \right) + O(\varepsilon^2)$$

EXAMPLE

$$y'y'' - \varepsilon y^2 = 0, \quad y(0) = 0 \quad y'(0) = 1$$

Assume

$$y(t, \varepsilon) = y_0(t) + \varepsilon y_1(t) + \varepsilon^2 y_2(t) + \dots$$

Initial condition expansions are clear

$$y_0(0) = 0 \quad y_0'(0) = 1$$

$$y_1(0) = 0 \quad y_1'(0) = 0$$

Expand differential equation

$$(y_0' + \varepsilon y_1' + \dots)(y_0'' + \varepsilon y_1'' + \dots) - \varepsilon (y_0 + \varepsilon y_1 + \dots)^2 = 0$$

Collecting like powers

$$O(1) \quad y_0' y_0'' = 0$$

$$O(\varepsilon) \quad y_0' y_1'' + y_0'' y_1 = y_0^2$$

are nonlinear equations

Solving  $O(1)$  problem

$$y_0' y_0'' = 0$$

$$y_0(0) = 0 \quad y_0'(0) = 1$$

Note that

$$\frac{d}{dt} (y_0'^2) = 0$$

Thus  $y_0'$  is constant and

$$y_0(t) = At + B$$

Applying initial conditions

$$\begin{aligned} y_0'(0) &= A = 1 \\ y_0(0) &= B = 0 \end{aligned}$$

Conclude

$$\boxed{y_0(t) = t}$$

Solving  $O(\varepsilon)$  problem

Using  $y_0(t) = t$  in  $O(\varepsilon)$  eqns we arrive at

$$y_1'' = t^2 \quad y_1(0) = 0 \quad y_1'(0) = 0$$

whose general solution is

$$y_1(t) = \frac{1}{12} t^4 + At + B$$

for constants  $A, B$ .

$$y_1(0) = B = 0$$

$$y_1'(0) = A = 0$$

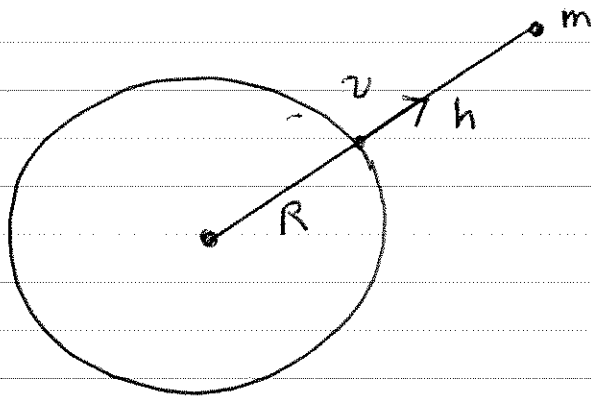
Conclude

$$y_1(t) = \frac{1}{12} t^4$$

and

$$y(t, \varepsilon) = t + \frac{1}{2} \varepsilon t^4 + O(\varepsilon^2)$$

## Motivational Example - Projectile Problem



A projectile of mass  $m$  is launched from earth surface at a velocity  $v$ .

Earth has radius  $R$  and mass  $M$ .  $G$  is the universal gravitational constant.

The height  $h(\pi)$  of the projectile is found using Newton's law of gravity

$$(1) \quad m \frac{d^2 h}{d\pi^2} = - \frac{GMm}{(h+R)^2}, \quad h(0) = 0, \quad h'(0) = v$$

Convert this dimensional IVP into a dimensionless one

$$y = \frac{h}{H^*} \quad t = \frac{\pi}{\pi^*}$$

yields

$$\frac{d^2 y}{dt^2} = - \left( \frac{GM\pi^{*2}}{H^*R^2} \right) \frac{1}{\left(1 + \frac{H^*}{R}y\right)^2} \quad y'(0) = \frac{v\pi^*}{H^*}$$

choose  $H^*$  and  $\pi^*$  so these both equal 1

yields

$$\pi^* = \frac{vR^2}{GM} = \frac{v}{g} \quad H^* = \frac{v^2 R^2}{GM}$$

Dimensionless problem - (what's small?)

$$(2) \quad y'' = -\frac{1}{(1+\varepsilon y)^2} \quad y(0) = 0, \quad y'(0) = 1$$

where the sole parameter is

$$(3) \quad \varepsilon = \frac{v^2 R}{GM} \ll 1$$

To see why  $\varepsilon$  might be small note that the gravitational constant  $g$  at the surface is

$$g = \frac{GM}{R^2} = 9.8 \text{ m/sec}^2$$

Then  $\varepsilon$  can be written

$$(4) \quad \varepsilon = \frac{v^2}{gR} \ll 1$$

owing to the fact that the radius  $R$  of the earth is large

$$R \approx 6.4 \times 10^6 \text{ m} \quad g = 9.8 \text{ m/sec}^2$$

A couple of examples

$$\varepsilon \approx 0.00016$$

$$v = 100 \text{ m/sec}$$

$$\varepsilon \approx 0.017$$

$$v = 1000 \text{ m/sec} \\ (\text{Mach } 3.0)$$

## Regular series approximation

$$(1) \quad y'' = - \frac{1}{(1 + \epsilon y)^2}, \quad 0 < \epsilon \ll 1$$

$$(2) \quad y(0) = 0$$

$$(3) \quad y'(0) = 1$$

Assume

$$y(t, \epsilon) = y_0(t) + \epsilon y_1(t) + \epsilon^2 y_2(t) + O(\epsilon^3)$$

Expand initial conditions

$$y(0, \epsilon) = y_0(0) + \epsilon y_1(0) + \dots = 0$$

$$y'(0, \epsilon) = y_0'(0) + \epsilon y_1'(0) + \dots = 1$$

yields

$$y_0(0) = 0 \quad y_0'(0) = 1$$

$$y_1(0) = 0 \quad y_1'(0) = 0$$

To expand the differential equation we use the Binomial Theorem

$$\frac{1}{(1 + \epsilon y)^2} = (1 + \epsilon y)^{-2} = 1 - 2\epsilon y + 3\epsilon^2 y^2 + \dots$$

Recall Binomial Thm

$$(1 + z)^p = 1 + pz + \frac{1}{2!} p(p-1)z^2 + \dots$$

if  $|z| < 1$

$$\begin{aligned} \frac{1}{(1+\varepsilon y)^2} &= 1 - 2\varepsilon(y_0 + \varepsilon y_1 + \dots) + 3\varepsilon^2(y_0 + \varepsilon y_1 + \dots)^2 + \dots \\ &= 1 - 2\varepsilon y_0 + (3y_0^2 - 2y_1)\varepsilon^2 + O(\varepsilon^2) \end{aligned}$$

Since

$$y'' = y_0'' + \varepsilon y_1'' + \varepsilon^2 y_2'' + \dots$$

the expansions above yield (by equating like powers of  $\varepsilon$ ) the following series of initial value problems

$$\begin{array}{lll} (4) & y_0'' = -1 & y_0(0) = 0 \quad y_0'(0) = 1 \\ (5) & y_1'' = 2y_0 & y_1(0) = 0 \quad y_1'(0) = 0 \\ (6) & y_2'' = -3y_0^2 + 2y_1 & y_2(0) = 0 \quad y_2'(0) = 0 \end{array}$$

whose solutions are

$$y_0(t) = -\frac{1}{2}t^2 + t$$

$$y_1(t) = -\frac{1}{12}t^3(t-4)$$

$$y_2(t) = -\frac{1}{360}t^4(11t^2 - 66t + 90)$$



## Summary and application

$$y(t, \varepsilon) = (t - \frac{1}{2}t^2) - \frac{1}{12}\varepsilon t^3(t-4) + O(\varepsilon^2)$$

can convert back to dimensional form using

$$y = \frac{g}{v^2} h \quad t = \frac{g}{v} \tau$$

## Maximal height correction

Seek an approximation to the time the projectile takes to achieve its maximal height. This occurs when its velocity is zero:

$$y' = f(t, \varepsilon) = 1 - t + \varepsilon(t^2 - \frac{1}{3}t^3) + \dots = 0$$

Let

$$t = \tau(\varepsilon) = \tau_0 + \varepsilon\tau_1 + O(\varepsilon^2)$$

Substitution and collection of powers yields

$$(1 - \tau_0) + \varepsilon(\tau_0^2 - \frac{1}{3}\tau_0^3 - \tau_1) + O(\varepsilon^2) = 0$$

From which  $\tau_0 = 1$ ,  $\tau_1 = \frac{2}{3}$  and

$$\tau = 1 + \frac{2}{3}\varepsilon + O(\varepsilon^2)$$

represents a slightly increased time to maximum height - as one would expect since gravity weakens the higher one goes.