

Relaxation oscillators

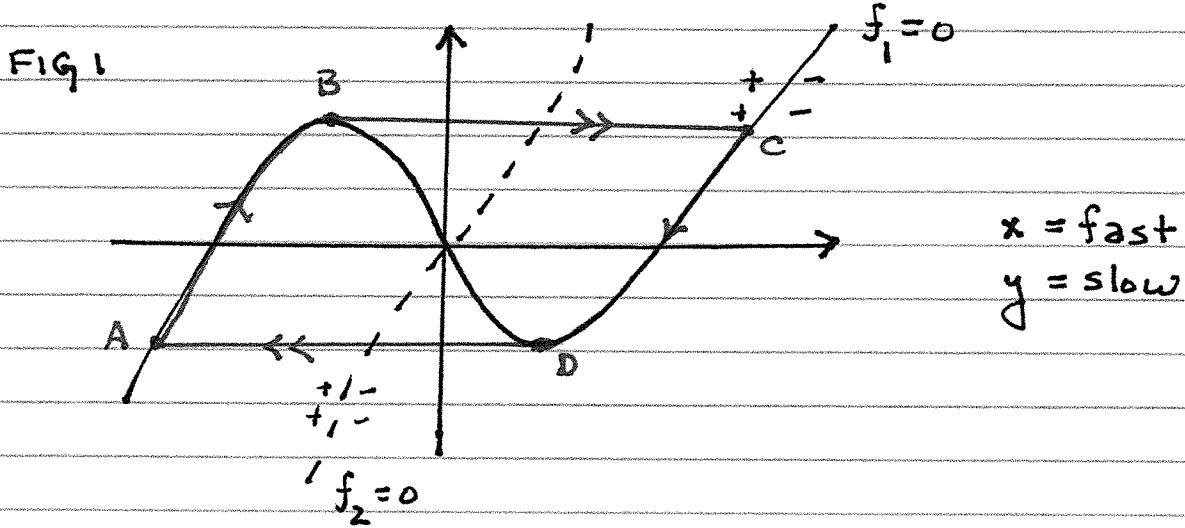
Is a periodic orbit where one dependent variable varies much quicker than the other. Eqns are

$$(1) \quad \dot{x} = f_1(x, y)$$

$$(2) \quad \dot{y} = \varepsilon f_2(x, y) \quad 0 < \varepsilon \ll 1$$

where ε is a small positive number.

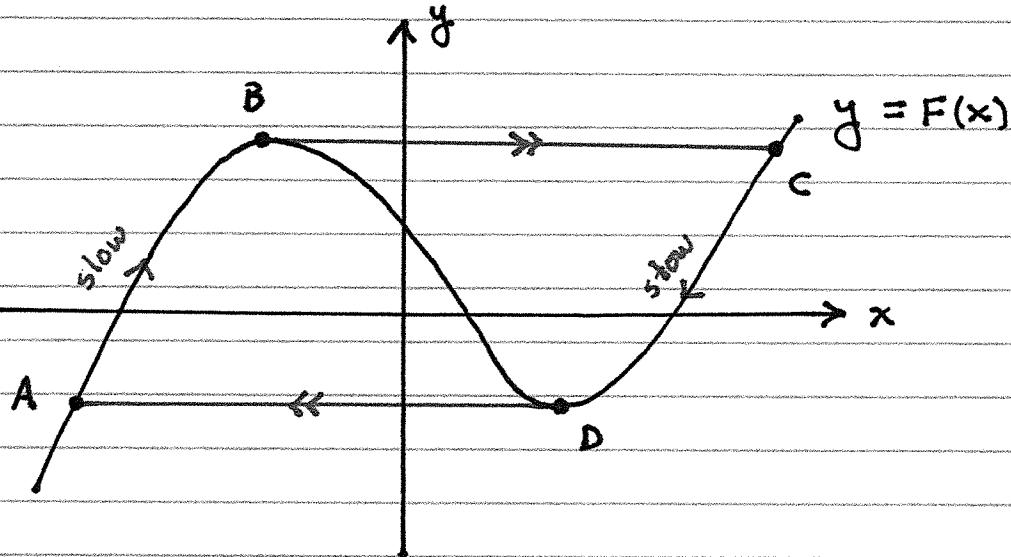
Qualitatively the nullclines have the assumed shapes



Can prove (Poincaré) periodic orbit near ABCD exists and that its period T is

$$T \approx T_{AB} + T_{CD}$$

Computing the approximate period



We assume that on the x -nullcline $f_1(x, y) = 0$ implies $y = F(x)$ for some function $F(x)$.

We use the notational convention

$$\alpha = (x_\alpha, y_\alpha) \quad \alpha = A, B, C, D$$

Introduce a slow time τ

$$\tau = \varepsilon t$$

so the differential equations are

$$(3) \quad \varepsilon \frac{dx}{d\tau} = f_1(x, y)$$

$$(4) \quad \frac{dy}{d\tau} = f_2(x, y)$$

Note that when $\varepsilon = 0$ in (3) we have

$$f_1(x, y) = 0$$

or that the solution lies on the $f = 0$ nullcline as depicted in the figure along AB and CD. BC and DA are rapid transitions. Of course in reality ε is only small hence trajectories are just near that nullcline.

Seek a leading order approximation by the expansions

$$x(\tau) = x_0(\tau) + \varepsilon x_1(\tau) + \dots$$

$$y(\tau) = y_0(\tau) + \varepsilon y_1(\tau) + \dots$$

The leading equations are:

$$(5) \quad 0 = f_1(x_0, y_0) \quad \left. \begin{array}{l} \\ \end{array} \right\} \text{"leading order slow subsystem"}$$

$$(6) \quad \frac{dy_0}{d\tau} = f_2(x_0, y_0)$$

But since (5) $\Rightarrow y_0 = F(x_0)$ so

$$(7) \quad \frac{dy_0}{d\tau} = H(x_0) \equiv f_2(x_0, F(x_0))$$

But again since $y_0 = F(x_0)$ we have

$$\frac{dy_0}{dt} = F'(x_0) \frac{dx_0}{dt}$$

Using this in (7) we obtain a decoupled differential equation for x_0

$$(8) \quad \frac{dx_0}{dt} = \frac{H(x_0)}{F'(x_0)}$$

So, from this we deduce

$$T_{AB} = \int_{x_A}^{x_B} \frac{F'(x)}{H(x)} dx$$

and a similar expression for $T_{CD} \Rightarrow$

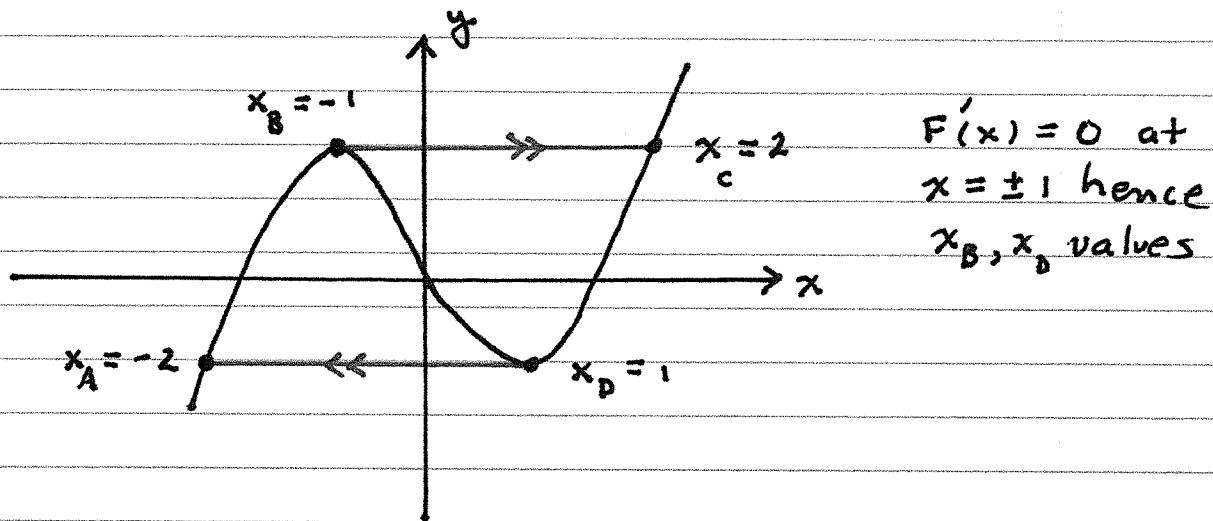
$$T \approx \int_{x_A}^{x_B} \frac{F'(x)}{H(x)} dx + \int_{x_C}^{x_D} \frac{F'(x)}{H(x)} dx$$

EXAMPLE

$$\dot{x} = y - F(x)$$

$$\dot{y} = -\epsilon x$$

$$F(x) \equiv \frac{1}{3}x^3 - x$$



$F'(x) = 0$ at
 $x = \pm 1$ hence
 x_B, x_D values

Leading order slow subsystem $T = \epsilon t$

$$(1) \quad y = F(x)$$

$$(2) \quad y' = -x$$

from which we find $y' = F'(x)x' \Rightarrow$

$$\frac{dx}{dt} = -\frac{x}{F'(x)}$$

$$F'(x) = x^2 - 1$$

Separate and integrate

$$T_{AB} = \int_{-2}^{-1} -\frac{F'(x)}{x} dx = 3 - 2\ln 2$$

By symmetry $T_{CD} = T_{AB}$ so total period is

$$T \approx 2T_{AB} = 6 - 4\ln 2$$