

INTEGRAL EQUATIONS - Classification

FREDHOLM INTEGRAL EQN ($\phi = \text{unknown}$)

$$\int_a^b K(x, y) \phi(y) dy = f(x) \quad (\text{1st kind})$$

$$\phi(x) - \lambda \int_a^b K(x, y) \phi(y) dy = f(x) \quad (\text{2nd kind})$$

VOLTERRA INTEGRAL EQN ($\phi = \text{unknown}$)

$$\int_a^x K(x, y) \phi(y) dy = f(x) \quad (\text{1st kind})$$

$$\phi(x) - \lambda \int_a^x K(x, y) \phi(y) dy = f(x) \quad (\text{2nd kind})$$

LINEAR INTEGRAL EQNS (EX)

$$K(x, \phi) \equiv K_1(\phi(x)) + \int_a^b K_2(x, y, \phi(y)) dy$$

$$K(x, \phi_1 + \phi_2) = K(x, \phi_1) + K(x, \phi_2)$$

$$K(x, \alpha \phi_1) = \alpha K(x, \phi_1)$$

CONNECTION BETWEEN FREDHOLM/VOLTERRA INT. EQNS

FOR VOLTERRA, $K(x, y) = 0$ for $y > x$

SOLVED FREDHOLM INTEGRAL EQUATION

$$u(x) - \lambda \int_0^1 e^{x-y} u(y) dy = f(x) \quad (1)$$

unknown here is $u(x)$.

Soln 1 Define $\eta = \int_0^1 e^{-y} u(y) dy$ then (1) becomes

$$u(x) - \lambda e^x \eta = f(x)$$

$$u(x) = f(x) + \eta \lambda e^x \quad (2)$$

where η is unknown. Using (2) in defn of η we find

$$\eta = \frac{1}{1-\lambda} \int_0^1 e^{-y} f(y) dy, \quad \lambda \neq 1$$

Thus,

$$u(x) = f(x) + \frac{\lambda}{1-\lambda} \int_0^1 e^{x-y} f(y) dy \quad \square$$

Remark: Defining the operator $Ku = \int_0^1 e^{x-y} u(y) dy$

$$(I - \lambda K)u = f$$

Then the solution is

$$u = \frac{1}{1-\lambda} (I + \lambda K) f$$

Remark: when $\lambda = 1$, note that $v(x) = Ae^x$ has

$$v(x) - \int_0^1 Ae^x dy = 0 \quad \forall A \in \mathbb{R}$$

so that $u(x) = Ae^x + f(x)$ is a (nonunique) soln family.

SOLVED VOLTERA EQUATION VIA LAPLACE TRANSFORM

$$(1) \quad u(t) = 1 + \int_0^t (t-\tau) u(\tau) d\tau \quad t > 0$$

Has the form

$$u(t) = f(t) + \int_0^t k(t-\tau) u(\tau) d\tau$$

Recall facts from Laplace transform theory

$$F(s) = \mathcal{L}(f)(s) \equiv \int_0^{\infty} e^{-st} f(t) dt$$

Given convolution defn we have

$$(f * g)(t) \equiv \int_0^t f(t-\tau) g(\tau) d\tau$$

$$\mathcal{L}(f * g)(s) = F(s) G(s) = \mathcal{L}(f) \cdot \mathcal{L}(g)$$

Taking the \mathcal{L} -transform of (1), $U(s) \equiv \mathcal{L}(u)(s)$

$$U(s) = \mathcal{L}(1) + \mathcal{L}(t) \mathcal{L}(u)$$

$$U(s) = \frac{1}{s} + \frac{1}{s^2} U(s)$$

$$U(s) = \frac{s}{s^2 - 1} \quad \text{Transform of soln.}$$

Use Inversion formula (or other means)

$$u(t) = \frac{1}{2\pi i} \int_{\gamma-i\infty}^{\gamma+i\infty} e^{st} U(s) ds = \sum_{s=\pm 1} \text{Res } \tilde{f}(s, t) = \cosh t$$

REMARK General coupled linear eqns can be solved same way.

Some simple conversions to Fredholm integral eqns

Claim

Solution of

$$u''(x) = f(x)$$

$$u(0) = u(1) = 0$$

$$u(x) = \int_0^1 k(x,y) f(y) dy$$

$$k(x,y) = \begin{cases} y(x-1) & 0 \leq y < x \leq 1 \\ x(y-1) & 0 \leq x < y \leq 1 \end{cases}$$

First note $k(0,y) = k(1,y) = 0$ so that $u(0) = u(1) = 0$.

$$u(x) = \int_0^x y(x-1) f(y) dy + \int_x^1 x(y-1) f(y) dy$$

$$u'(x) = \int_0^x y f(y) dy + \int_x^1 (y-1) f(y) dy$$

$$u''(x) = x f(x) - (x-1) f(x)$$

$$u''(x) = f(x) \quad \square$$

EXAMPLE Convert to a Fredholm integral equation.

$$u'' + \lambda u = f(x)$$

$$u(0) = u(1) = 0$$

Set $F(x) = f(x) - \lambda u(x)$ and apply claim

$$u(x) + \lambda \int_0^1 k(x,y) u(y) dy = g(x)$$

where

$$g(x) = \int_0^1 k(x,y) f(y) dy \quad \left(\begin{array}{l} \text{Comment on nonuniqueness} \\ \text{if } \lambda_n = n\pi \end{array} \right)$$

EXAMPLE

$$\begin{aligned} u'' + a(x)u' + b(x)u &= f(x) & x \in (0, 1) \\ u(0) = u(1) &= 0 \end{aligned}$$

Set

$$(1) \quad u(x) = \phi(x) \psi(x) \quad \psi(x) = \exp\left(-\frac{1}{2} \int_0^x a(s) ds\right)$$

to eliminate first derivatives. Since $\psi > 0$ must have

$$(2) \quad \phi(0) = \phi(1) = 0$$

Substitute (1) into ODE yields

$$(3) \quad \phi'' + q(x)\phi = \psi^{-1}f$$

$$(4) \quad q(x) = b(x) - \frac{1}{4}a^2(x) - \frac{1}{2}a'(x)$$

Then, using claim

$$(5) \quad \phi(x) + \int_0^1 \hat{k}(x, y) \phi(y) dy = \hat{f}(x)$$

where

$$(6) \quad \hat{f}(x) = \int_0^1 k(x, y) \psi^{-1}(y) f(y) dy$$

$$(7) \quad \hat{k}(x, y) = q(y) k(x, y)$$

Conversion of IVP to Volterra Eqns (Linear eqn case)

$$(1) \quad L[u] = u^{(n)}(x) + a_1(x)u^{(n-1)}(x) + \dots + a_n(x)u = F(x)$$

$$(2) \quad u(0) = c_0 \quad u'(0) = c_1, \dots, \quad u^{(n-1)}(0) = c_{n-1}$$

Define

$$D^n u = \frac{d^n u}{dx^n} = \phi(x)$$

and for any $\psi(x)$

$$(3) \quad D^{-1}\psi \equiv \int_0^x \psi(y) dy$$

Using (Laplace transforms or) integration by parts

$$(4) \quad D^{-n}\psi = \frac{1}{(n-1)!} \int_0^x (x-y)^{n-1} \psi(y) dy$$

Proof for $n=2$

$$\begin{aligned} D^{-2}\psi &= D^{-1}(D^{-1}\psi) = \int_0^x \overbrace{\left(\int_0^s \phi(y) dy \right)}^{u(s)} \overbrace{ds}^{dv(s)} \\ &= s \int_0^s \phi(y) dy \Big|_{s=0}^{s=x} - \int_0^x s \phi(s) ds \\ &= x \int_0^x \phi(y) dy - \int_0^x s \phi(s) ds \\ &= \int_0^x (x-y) \phi(y) dy \quad \square \end{aligned}$$

(See back for convolution method?)

Now define the n -term Taylor series about $x=0$ of u

$$U(x) = \sum_{k=0}^n \frac{u^{(k)}(0)}{k!} x^k \quad \text{where } u^{(0)}(x) \equiv u(x)$$

Using initial conditions

$$U(x) = \sum_{k=0}^n \frac{c_k}{k!} x^k$$

Now observe

$$(5) \quad \begin{cases} \frac{d^{n-1} u}{dx^{n-1}} = c_{n-1} + D^{-1} \phi & = U^{(n-1)}(x) + D^{-1} \phi \\ \frac{d^{n-2} u}{dx^{n-2}} = c_{n-1} x + c_{n-2} + D^{-2} \phi & = U^{(n-2)}(x) + D^{-2} \phi \\ \vdots & \vdots \\ u = \frac{c_{n-1}}{(n-1)!} x^{n-1} + \dots + c_1 x + c_0 + D^{-n} \phi & = U(x) + D^{-n} \phi \end{cases}$$

Using (5), (4) in (1) we find (1) can be written

$$\phi(x) + \int_0^x K(x,y) \phi(y) dy = f(x), \quad \phi = \frac{d^n u}{dx^n}$$

where

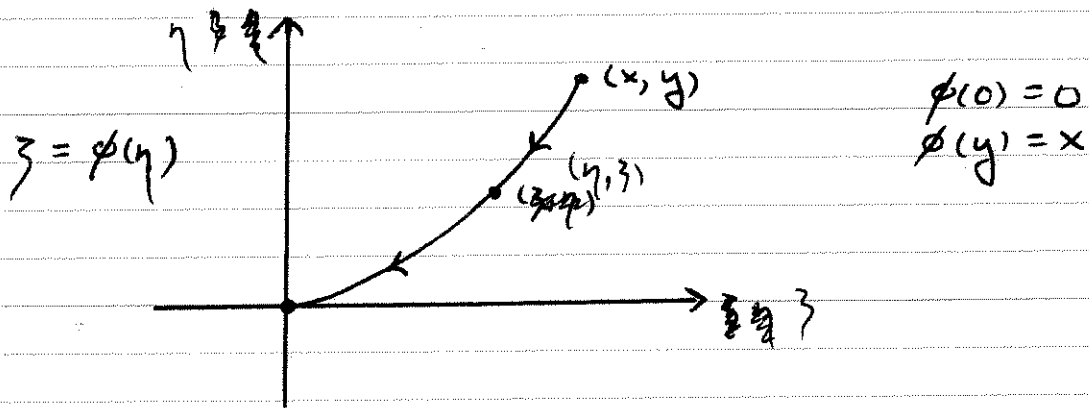
$$K(x,y) = \sum_{j=1}^n a_j(x) \frac{(x-y)^{j-1}}{(j-1)!}$$

$$f(x) = F(x) - \sum_{j=1}^n a_j(x) U^{(n-j)}(x)$$

ABEL'S EQUATION

A particle of mass m initially at rest at (x, y) moves on curve $z = \phi(\eta)$ in the (z, η) -plane through the origin.

If $T(y)$ is the time taken to traverse the path, what is $\phi(\eta)$?



Total energy is $E = 0$. Conservation of energy

$$\frac{1}{2} m \left(\frac{ds}{dt} \right)^2 - mg(y - \eta) = 0$$

where $s = \text{arc length}$, $\frac{ds}{dt} = \text{speed}$, implies

$$(1) \quad \frac{ds}{dt} = \sqrt{2g(y - \eta)}$$

Basic calculus

$$(2) \quad ds = \sqrt{1 + (\phi'(\eta))^2} d\eta$$

Combining (1)-(2) gives

$$\frac{\sqrt{1 + (\phi'(\eta))^2} d\eta}{\sqrt{2g(y-\eta)}} = dt$$

Since traversal time $T(y)$ is known

$$\int_0^y \frac{\sqrt{1 + (\phi'(\eta))^2} d\eta}{\sqrt{2g(y-\eta)}} = T(y)$$

Defining $\psi(\eta) = (2g)^{-1/2} (1 + (\phi'(\eta))^2)^{1/2}$ we get

$$\int_0^y \frac{\psi(\eta) d\eta}{\sqrt{y-\eta}} = T(y)$$

Abel's Integral Equation.

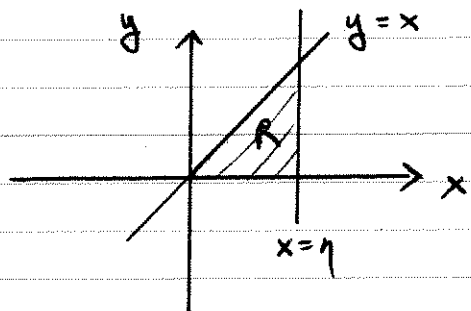
Knowing the solution $\psi(\eta)$, one can solve for $\phi(\eta)$ given initial condition $\phi(0) = 0$.

SOLUTION OF ABEL'S INTEGRAL EQN. (METHOD ONE)

$$T(x) = \int_0^x \frac{u(y) dy}{\sqrt{x-y}}$$

Employ a trick. Multiply by $\frac{dx}{\sqrt{\eta-x}}$ and integrate from $x=0$ to $x=\eta$

$$\int_0^{\eta} \frac{T(x) dx}{\sqrt{\eta-x}} = \int_0^{\eta} \left(\int_0^x \frac{u(y)}{\sqrt{\eta-x} \sqrt{x-y}} dy \right) dx$$



Integration region of RHS

$$\left(\begin{array}{l} y < x < \eta \\ 0 < y < \eta \end{array} \right)$$

Interchange order of limits

$$\int_0^{\eta} \frac{T(x) dx}{\sqrt{\eta-x}} = \int_0^{\eta} u(y) \left(\int_y^{\eta} \frac{dx}{\sqrt{\eta-x} \sqrt{x-y}} \right) dy$$

Direct integration: equals π

$$\int_0^{\eta} \frac{T(x) dx}{\sqrt{\eta-x}} = \pi \int_0^{\eta} u(y) dy$$

Differentiating this result

$$u(\eta) = \frac{1}{\pi} \frac{d}{d\eta} \int_0^{\eta} \frac{T(x)}{\sqrt{\eta-x}} dx$$

REMARK: Laplace Transforms + Convolution gives another technique.

Solutions for separable kernels

$$Ku = \int_a^b \underbrace{\left(\sum_{i=1}^n \phi_i(x) \psi_i(y) \right)}_{k(x,y)} u(y) dy$$

Int eqn: Fredholm and kind

$$u(x) = f(x) + \lambda (Ku)(x)$$

Define $\alpha_j = \langle \psi_j, u \rangle$ then ($L^2[a,b]$ inn. prod)

$$u(x) = f(x) + \lambda \sum_{i=1}^n \alpha_i \phi_i(x) \quad (1)$$

$$\langle \psi_j, u \rangle = \langle \psi_j, f \rangle + \lambda \sum_{i=1}^n \alpha_i \langle \psi_j, \phi_i \rangle$$

$$\alpha_j = \langle \psi_j, f \rangle + \lambda \sum_{i=1}^n \langle \psi_j, \phi_i \rangle \alpha_i$$

Let $\alpha = (\alpha_1, \dots, \alpha_n)^T$, $\hat{f} = (\langle \psi_1, f \rangle, \dots, \langle \psi_n, f \rangle)$ and $A \in \mathbb{R}^{n \times n}$, $[A]_{ij} = \langle \psi_j, \phi_i \rangle$ gives

$$\alpha = \hat{f} + \lambda A \alpha$$

$$(I - \lambda A) \alpha = \hat{f}$$

If $(I - \lambda A)^{-1}$ exists, (1) gives the soln with $\alpha = (I - \lambda A)^{-1} \hat{f}$.

Sep Int Eqns \Leftrightarrow Matrix Eqns

EXAMPLE

$$u(x) = 1 + \lambda \int_0^1 (xy^2 + x^2y) u(y) dy \quad (1)$$

where $f(x) = 1$,

$$k(x, y) = xy^2 + x^2y = \phi_1(x)\psi_1(y) + \phi_2(x)\psi_2(y)$$

Define

$$\alpha_1 = \langle \psi_1, u \rangle = \int_0^1 y^2 u(y) dy$$

$$\alpha_2 = \langle \psi_2, u \rangle = \int_0^1 y u(y) dy$$

then (1) is

$$u(x) = 1 + \lambda (\alpha_1 \phi_1(x) + \alpha_2 \phi_2(x)) \quad (2)$$

Define

$$A = \begin{bmatrix} \langle \psi_1, \phi_1 \rangle & \langle \psi_1, \phi_2 \rangle \\ \langle \psi_2, \phi_1 \rangle & \langle \psi_2, \phi_2 \rangle \end{bmatrix} = \begin{bmatrix} \frac{1}{4} & \frac{1}{5} \\ \frac{1}{3} & \frac{1}{4} \end{bmatrix}$$

$$F = \begin{bmatrix} \langle f, \psi_1 \rangle \\ \langle f, \psi_2 \rangle \end{bmatrix} = \begin{bmatrix} \frac{1}{3} \\ \frac{1}{2} \end{bmatrix}$$

Then $\langle \psi_j, \cdot \rangle$ on (2) gives

$$(I - \lambda A) \alpha = F \quad (3)$$

where

$$\det(I - \lambda A) = 1 - \frac{1}{2}\lambda - \frac{1}{240}\lambda^2$$

Solving (3) for $\alpha = (\alpha_1, \alpha_2)^T$ and substituting into (2)

$$u(x) = 1 + \frac{2}{3} \frac{\lambda x [5x(\lambda - 36) - 6(\lambda + 20)]}{\lambda^2 + 120\lambda - 240}$$

is the solution

Eigenvalue problem

$$u(x) = \lambda \int_0^1 k(x,y) u(y) dy = \lambda Ku$$

re-written

$$Ku = \frac{1}{\lambda} u$$

thus $\frac{1}{\lambda}$ is e-value of K . For $f=0$ we find e-val
eqn is

$$\det(I - \lambda A) = 0$$

$$1 - \frac{1}{2}\lambda - \frac{1}{240}\lambda^2 = 0$$

$$\lambda_{\pm} = -60 \pm 16\sqrt{15}$$

Eigenfns

$$u(x) = \alpha_{\pm}^T (\phi_1(x), \phi_2(x)) \quad \alpha_{\pm} \in N(I - \lambda_{\pm} A)$$

EIGENVALUE EXAMPLE

Let $D = \{(x, y) : x^2 + y^2 \leq 1\}$ and define

$$Ku = \iint_D (x^2 + y^2) u(x, y) dA.$$

Compute eigenvalue and normalized e-fn. ($L^2(D)$)

$$Ku = \lambda u = \alpha_1, \quad \alpha_1 \in \mathbb{R}$$

Thus $u = \alpha_1 / \lambda$ and ($\lambda \neq 0$)

$$Ku = \int_0^{2\pi} \int_0^1 r^2 \cdot \left(\frac{\alpha_1}{\lambda}\right) r dr d\theta$$

$$Ku = 2\pi \cdot \frac{1}{4} \cdot \frac{\alpha_1}{\lambda} = \alpha_1$$

From which

$$\lambda = \frac{\pi}{2}$$

Since $u = \text{const} = \bar{u}$, need

$$\|u\|^2 = \iint_D \bar{u}^2 dA = \pi \bar{u}^2 = 1$$

So

$$(\lambda, u) = \left(\frac{\pi}{2}, \frac{1}{\sqrt{\pi}}\right)$$

For $\lambda = 0$ eval, note $\phi_n(r, \theta) \in \{\sin n\theta, \cos n\theta\}$

$$K\phi_n = 0.$$

EXAMPLE

$$Ku \equiv \int_a^b k(x,y) u(y) dy = \sum_{i=1}^N \phi_i(x) \int_a^b \psi_i(y) u(y) dy$$

The range $R(K) \in \text{span} \{ \phi_i \}$ thus $\dim R(K) = N$.

If ϕ_i independent, $N(K) = \{0\}$ and $\lambda = 0$ not e-value.

$$\alpha_i \phi_i(x) = 0 \quad \forall x \in (a,b)$$

Eigenvalue problem

$$Ku = \lambda u$$

$$\alpha_i \phi_i = \lambda u$$

$$(A - \lambda I) \alpha = 0$$

Thus

$$\det(A - \lambda I) = 0 \quad \text{e-value eqn}$$

$$u = \frac{1}{\lambda} \sum \alpha_i \phi_i(x) \quad \text{e-fns}$$

where $\alpha \in N(A - \lambda I)$.

Systems of integral eqns (Fred and kind)

$$\vec{u} + \lambda K \vec{u} = \vec{f}$$

where $\vec{u} = (u_1, \dots, u_n)$, $\vec{f} = (f_1, \dots, f_n)$ and

$$[K \vec{u}]_i = \int_a^b \sum_{j=1}^n k_{ij}(x,y) u_j(y) dy.$$

Solution via complete sets

let $\{\phi_n\}$ be a complete orthonormal set for $L^2[a,b]$.

$$u - \lambda K u = f$$

let $u = \sum_{i=1}^n u_i \phi_i(x)$, $f = \sum_{i=1}^n f_i \phi_i(x)$. Set

$$(I) \quad Ku = \sum_{i=1}^n d_i K \phi_i$$

$$\sum (\phi_i - \lambda K \phi_i) u_i = \sum f_i \phi_i$$

$$\langle \phi_j, \phi_j - \lambda K \phi_j \rangle u_j = f_j$$

determines u_j for solution. But

1) How do we know (I) is valid?

2) If $\frac{1}{\lambda}$ is an eval of K then some u_j (how many) cannot be found.

Goal of Chapter has to do with

$$(1) \quad u(x) + \int_a^b K(x,y) u(y) dy = f(x)$$

Suppose

$$Ku = \int_a^b k(x,y) u(y) dy, \quad K: L^2[a,b] \rightarrow L^2[a,b]$$

generates a basis for $H = L^2[a,b]$ via

$$K\phi_n = \lambda_n \phi_n, \quad \|\phi_n\| = 1$$

Then

$$u(x) = \sum_{n \geq 1} u_n \phi_n(x) \quad f(x) = \sum_{n \geq 1} f_n \phi_n(x)$$

in (1) gives

$$\sum u_n \phi_n + \sum u_n K\phi_n = \sum f_n \phi_n$$

$$\sum [(1 - \lambda_n) u_n - f_n] \phi_n = 0$$

Thus

$$u_n = \frac{f_n}{1 - \lambda_n}, \quad \lambda_n \neq 1$$

gives soln.