

## Calculus of Variations

Let  $V$  be some vector space of real valued functions. Now let  $J$  be a functional

$$J : A \rightarrow \mathbb{R}$$

where  $A \subset V$  is a subset (not necessarily subspace) of admissible functions  $u \in A$ .  
Goal is to minimize  $J$  over  $A$

$$\min_{u \in A} J(u)$$

In some texts if the min is attained by some  $\bar{u} \in A$  one writes

$$\bar{u} = \underset{u \in A}{\operatorname{argmin}} J(u)$$

Key theoretical issues

(1) Necessary conditions on solns

(2) Existence and uniqueness of solns

(3) Sufficient conditions

In practice the sets  $A$  may be very complex and defined by several constraints

Ex

### Point evaluation functional

$$J(y) = y'(1)^2$$

$$A = C[0, 2]$$

Ex

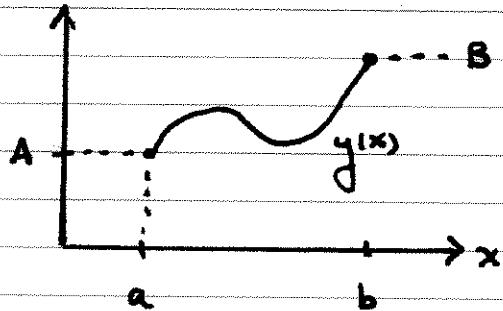
### Area functional

$$J(y) = \int_0^1 y(x) dx$$

A possible constraint might be  $y(x) \geq 0$

Ex

### Arclength functional



$$V = PC'[a, b]$$

piecewise  $C^1$

$$J(y) = \int_a^b \sqrt{1 + y'^2} dx$$

Here the admissible set

$$A = \{ y \in V : y(a) = A, y(b) = B \}$$

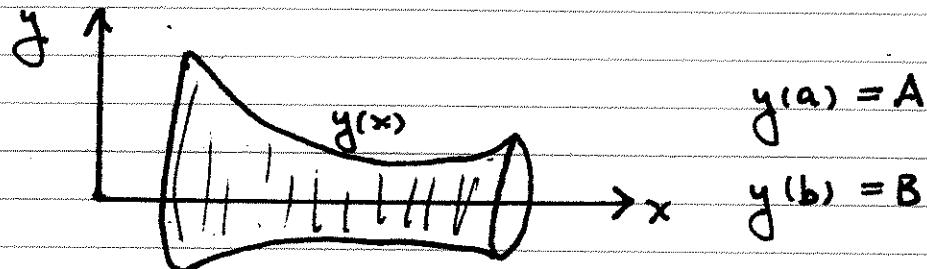
and is clearly not a subspace of  $V$ .

Also the minimizer is the straight line connecting endpoints

$$\bar{y}(x) = B + \frac{(B-A)}{(b-a)}(x-a)$$

EX

## Surface area of revolution



From elementary calculus the surface area of revolution is

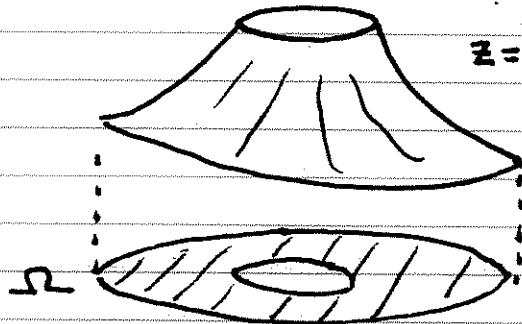
$$J(y) = \int_a^b 2\pi y \sqrt{1 + y'^2} dx$$

Surprisingly the answer is NOT part of a cone.

EX

## Minimum Area

$$V = C(n)$$



$$z = u(x, y)$$

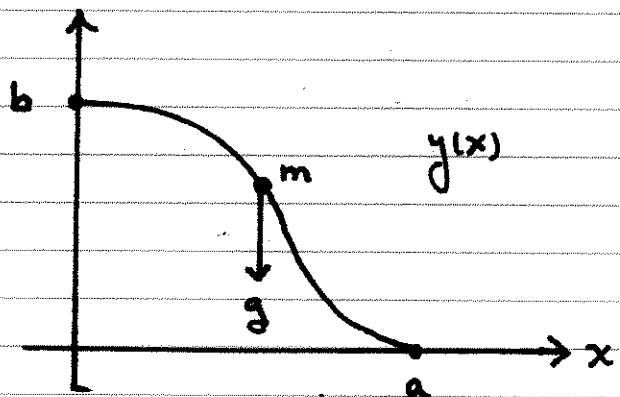
Given  $\Omega$  and the value of  $u$  on the boundary  $\partial\Omega$  we seek the surface of minimum area

$$J(u) = \int_{\Omega} \sqrt{1 + u_x^2 + u_y^2} dA$$

where

$$A = \{u \in C(\bar{\Omega}) : u|_{\partial\Omega} = g\}$$

Ex Brachistochrone Problem



A bead of mass  $m$  slides down a frictionless wire at speed  $v$  under the influence of gravity. Of all shapes  $y(x)$ , which minimizes the transit time  $T$ .

$$(1) \quad T = \int_0^T dt = \int_0^L \frac{dt}{ds} ds = \int_0^L \frac{1}{v} ds = \int_0^a \frac{\sqrt{1+y'^2}}{v} dx$$

Use conservation of energy to find  $v$

$$(2) \quad mg(b-y) = \frac{1}{2}mv^2 + mgy$$

Solve (2) for  $v$  to use in (1). Must minimize

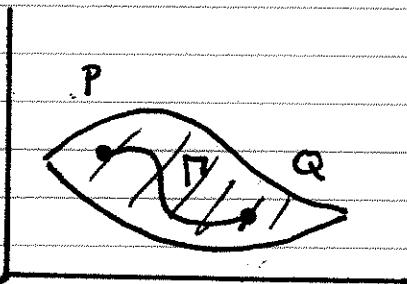
$$T(y) \equiv \int_0^a \sqrt{\frac{1+y'^2}{2g(b-y)}} dx$$

over admissible set

$$A = \{ y \in C^2[0, a] : y(0) = b, y(a) = 0 \}$$

EX

## Geodesics



$$g(x, y, z) = 0 \\ \text{surface}$$

of all curves  $\Gamma$  on  $g=0$ ,  
which has the shortest length.  
Such curves are called geodesics

$$\underline{x}(t) = (x(t), y(t), z(t)) \quad t \in [0, T]$$

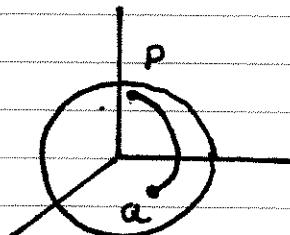
Seek to minimize

$$J(\underline{x}) = \int_0^T \sqrt{\dot{x}^2 + \dot{y}^2 + \dot{z}^2} dt$$

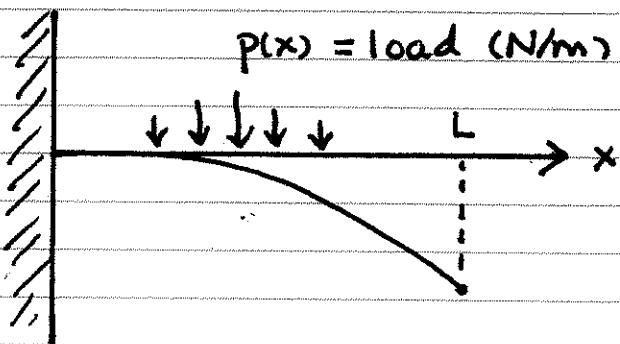
over an admissible set

$$A = \{ \underline{x} : \underline{x}(0) = P, \underline{x}(T) = Q, g(\underline{x}) = 0 \}$$

Note geodesics on spheres  
are "great circles"



Ex Beam deflection problem



Minimize potential energy

$$U(y) = \int_0^L [\frac{1}{2} \mu y''^2 - p(x)y] dx$$

where  $\mu$  = flexural rigidity (elastic component of pot. energy)

$$\min_{y \in A} U(y)$$

where

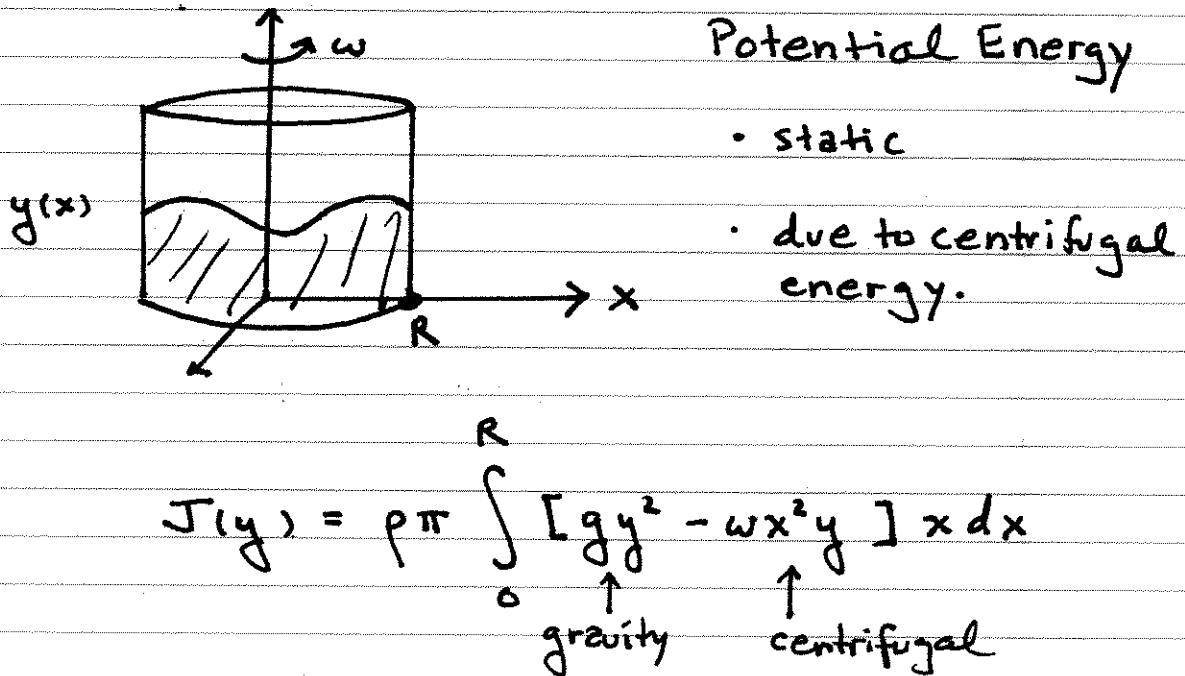
$$A = \{ y \in C^2[0, L] : y(0) = y'(0) = 0 \}$$

Note boundary conditions at  $x=L$  are "free".

Part of minimization procedure is to find the B.C. at  $x=L$ .

EX Rotating Fluid (Min Pot. Energy)

A fluid of density  $\rho$  in a cylinder of radius  $R$  is rotated with angular velocity  $\omega$



The integral arises from using the method of cylindrical shells to sum up contribution for P. Energy.

Has an "isoperimetric" constraint defining  $A$ .

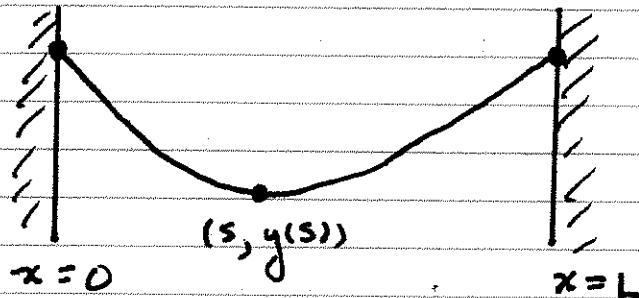
$$V(y) = 2\pi \int_0^R xy dx = \text{volume } V_0$$

Since volume of fluid is fixed

$$A = \{ y \in C[0, R] : V(y) = V_0 \}$$

Ex

## Hanging cable problem



Let  $s$  be arclength ( $s=0$  at  $x=0$ )  
and  $L_c$  be the length of  
the cable with  $L_c > L$ .

Shape is that which minimizes  
the potential energy

$$J(y) = \mu \int_0^{L_c} y(s) ds \quad \text{pot. energy}$$

where  $\mu = \text{weight/length}$ . Span fixed

$$K(y) = \int_0^L dx = \int_0^{L_c} \sqrt{1 + y'(s)^2} ds$$

follows from  $x'(s)^2 + y'(s)^2 = 1$  (unit tangent)

Seek  $y(s)$  to solve

$$\min_{y \in A} J(y)$$

where

$$A = \{ y : y(0) = y(L_c) = 0, K(y) = L \}$$

## Ex Classical Mechanics (nonrelativistic motion)

Let  $q \in \mathbb{R}^n$  be coordinates of a system of particles and  $\dot{q} = \frac{dq}{dt}$ .

$$J(q) = \int_0^t (T(q, \dot{q}) - U(q, \dot{q})) dt$$

is the "action" of the system when

$T(q, \dot{q})$  = total kinetic energy

$U(q, \dot{q})$  = total potential energy

In physics the principle of least action states the physically realized path is that which extremizes (minimizes)  $J$ .

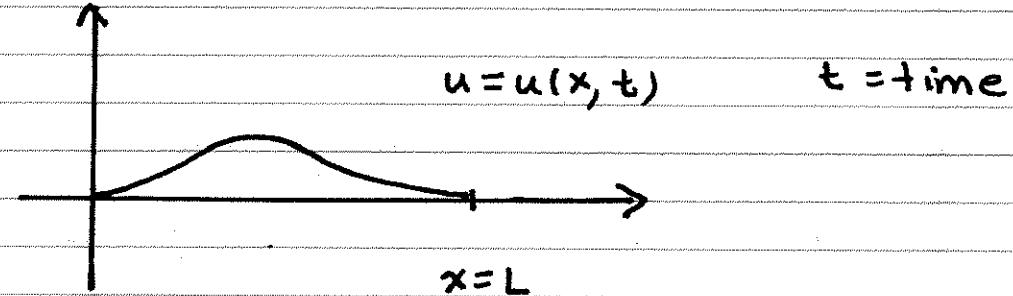
Note the action  $J$  is not the integral of the total energy,  $T + U$ .

Also, the statement is true for any "generalized" coordinates  $q$ , and leads to

1) Lagrangian mechanics

2) Hamiltonian mechanics

## Ex Taut string (least action continuum mech)



### Kinetic Energy

$$T = \frac{1}{2} \int_0^L \rho u_t^2 dx \quad \rho = \text{mass/length}$$

### Potential Energy

$$U = \int_0^L \mathcal{T}(x, t) \left( \sqrt{1 + u_x^2} - 1 \right) dx$$

where  $\mathcal{T}(x, t)$  = local tensile force.

Here  $U$  is the work done to stretch the string (only). Note  $U=0$  if  $u_x \equiv 0$ .

### Least Action

$$J(u) = \int_0^{t_L} \int_0^L (T - U) dx dt$$

extremize over  $u(x, t)$  satisfying

$$u(0, t) = u(L, t) = 0 \quad \text{B.C.}$$

$u(x, 0)$ ,  $u_t(x, 0)$  specified