

## Operator boundedness

Throughout  $L: H \rightarrow H$  is linear on Hilbert space  $H$ .

Defn  $L$  is a bounded operator if  $\exists M > 0$  s.t.

$$\|Lu\| \leq M \|u\| \quad \forall u \in H$$

Boundedness is akin to continuity:

Theorem (Stakgold for proof)

$$L \text{ continuous} \Leftrightarrow L \text{ bounded}$$

$$L \text{ continuous} \Leftrightarrow L(\lim_{n \rightarrow \infty} z_n) = \lim_{n \rightarrow \infty} Lz_n$$

EXAMPLE Let  $k(x,y) \in C[a,b]^2$

$$Lu = \int_a^b k(x,y)u(y)dy \quad \text{Bounded}$$

is a bounded operator on (all)  $H = L^2[a,b]$

Let  $v = Lu$

$$|v(x)|^2 = \left| \int_a^b k(x,y)u(y)dy \right|^2 \leq \|u\|^2 \int_a^b |k(x,y)|^2 dy$$

Hence

$$\|Lu\|^2 \leq \|u\|^2 \cdot \underbrace{\int_a^b \int_a^b |k(x,y)|^2 dy dx}_M$$

## Holder Inequalities

A set of inequalities for  $L_p(\Omega)$  spaces.  
For  $1 < p, q < \infty$

$$f \in L^p(\Omega) \Leftrightarrow \|f\|_p < \infty$$

where

$$\|f\|_p \equiv \left( \int_{\Omega} |f|^p dx \right)^{\frac{1}{p}}$$

Holder's inequality for

$$\frac{1}{p} + \frac{1}{q} = 1$$

is

$$|\langle f, g \rangle| \leq \|f\|_p \|g\|_q$$

with standard  $L^2$  inner product.

when  $p = q = 2$

$$|\langle f, g \rangle| \leq \|f\|_2 \|g\|_2$$

Longhand one can write (squaring above)

$$\left| \int_{\Omega} f(x)g(x) dx \right|^2 \leq \|f\|^2 \|g\|^2$$

EXAMPLE  $D(L) = BC[0, \infty)$  and continuous on  $[0, \infty)$

$$Lu \equiv \frac{1}{t} \int_0^t u(s) ds \quad (\text{average})$$

L'Hopital's rule implies  $u(0^+)$  defined!  
Bounded under sup norm  $\|u(t)\| = \sup_{t \geq 0} |u(t)|$

$$\left| \frac{1}{t} \int_0^t u(s) ds \right| \leq \frac{1}{t} \int_0^t |u(s)| ds \leq \|u\|$$

Hence  $L$  bounded on a Banach Space

EXAMPLE Derivative operator  $H = L^2[0, 2\pi]$

$$D(L) = C^1[0, 2\pi]$$

$$Lu \equiv \frac{du}{dx} \quad \text{unbounded}$$

Easy to show using  $u_n = \sin(nx)$ ,  $u_n' = n \cos(nx)$

$$\|u_n\|^2 = \int_0^{2\pi} \sin^2 nx dx = \pi$$

$$\|Lu_n\|^2 = \int_0^{2\pi} n^2 \cos^2 nx dx = n^2 \pi$$

Hence

$$\frac{\|Lu_n\|^2}{\|u_n\|^2} = n^2 \rightarrow \infty$$

and  $\nexists k$  s.t

$$\|Lu_n\| \leq k \|u_n\| \quad \forall n.$$

## Operators with nonclosed ranges $R(L)$

If  $R(L)$  is closed then  $H = R(L) \oplus R(L)^\perp$  needed for Fredholm alternative. If  $R(L)$  is not closed decomposition may not be possible. For  $H = \mathbb{R}^n$ ,  $Lu = Au$ ,  $A \in \mathbb{R}^{n \times n}$  the range is closed

EXAMPLE Even some bounded  $L$  have  $R(L) \neq \overline{R(L)}$

$$Lu \equiv \sum_{n=1}^{\infty} \frac{1}{n} \langle u, \phi_n \rangle \phi_n$$

where  $\{\phi_n\}$  orthonormal basis. Parseval equality  $\Rightarrow$

$$\|Lu\|^2 = \sum_{n \geq 1} \frac{1}{n^2} |\langle u, \phi_n \rangle|^2 \leq \sum_{n \geq 1} |\langle u, \phi_n \rangle|^2 = \|u\|^2$$

shows  $L$  bounded on  $H$ .

$$u_n \equiv \sum_{k=1}^n \phi_k \quad Lu_n = \sum_{k=1}^n \frac{1}{k} \phi_k, \quad \langle \phi_i, \phi_j \rangle = \delta_{ij}$$

Clearly  $Lu_n \rightarrow y \equiv \sum_{k=1}^{\infty} \frac{1}{k} \phi_k$ . In particular  $\|y\|^2 = \frac{\pi^2}{6}$ .

However,  $\nexists z \in H$  such that  $Lz = y$ . If there were

$$Lz = \sum_{n=1}^{\infty} \frac{1}{n} \langle z, \phi_n \rangle \phi_n = \sum_{n=1}^{\infty} \frac{1}{n} \phi_n$$

or  $\langle z, \phi_n \rangle = 1$  for all  $n$ . But then

$$z = \sum_{n=1}^{\infty} \phi_n \quad \text{diverges}$$

$L$ bounded
$Lu_n \rightarrow y \in H$
$\nexists z$ s.t. $Lz = y$

## Operator Domains

whether  $L$  is bounded or unbounded the domain  $D(L)$  must be a subspace of  $H$ .

If  $L$  is bounded it is almost always the case  $D(L) = H$ . If not its domain of definition can be extended.

Most unbounded  $L$  have 'dense' domains:

$L$  bounded

$$D(L) = H$$

$L$  unbounded

$D(L)$  dense in  $H$

EX  $Lu = Au$   $D(L) = \mathbb{R}^n$

EX  $Lu = \int_a^b K(x,y)u(y)dy$   $D(L) = L^2[a,b]$  \*\*

EX  $Lu = u''$  may have a domain

$$D(L) = \{u \in C^2[a,b] : u(a) = u(b) = 0\}$$

which is dense in  $L^2[a,b]$ .

This domain is associated with BVP

$$u'' = f \quad u(a) = u(b) = 0$$

\*\* assuming  $K \in H \times H$ ,  $H = L^2[a,b]$ .

## Adjoint Operators: Theory

Let  $L : D(L) \rightarrow H$ . We say  $L^* : D(L^*) \rightarrow H$  is the adjoint of  $L$  if

$$\langle Lu, v \rangle = \langle u, L^*v \rangle \quad \forall u \in D(L), \forall v \in D(L^*)$$

For  $H = \mathbb{R}^n$  the adjoint  $L^*$  of  $Lu \equiv Au$  is easily defined for

$$L^*u = A^T u$$

Moreover  $D(L) = D(L^*) = H$ .

### EXISTENCE FOR BOUNDED OPERATORS

If  $L$  is bounded operator on  $H$  then  $l : H \rightarrow \mathbb{C}$ ,  $l(u) \equiv \langle Lu, v \rangle$  is a bounded functional for each fixed  $v$ .

From theory on functionals (bounded) the Riesz representation theorem implies for each  $v \in H$ ,  $\exists g$  s.t.

$$\langle Lu, v \rangle = \langle u, g \rangle \quad \forall u \in H.$$

Since  $g$  depends uniquely on  $v$  we let

$$L^*v = g$$

namely an adjoint operator exists.

## EXISTENCE FOR UNBOUNDED OPERATORS

Let  $L: D(L) \rightarrow H$  be unbounded with a domain  $D(L)$  dense in  $H$ .

Seek to find  $g$  such that

$$(1) \quad \langle Lu, v \rangle = \langle u, g \rangle, \quad \forall u \in D(L)$$

Clearly  $(v, g) = (0, 0)$  works but seek  $g$ 's for  $v \neq 0$ .

Maybe (1) is true for  $g_1, g_2$ . Bad news.

But, were this the case

$$(3) \quad \langle u, g_1 - g_2 \rangle = 0 \quad \forall u \in D(L)$$

The denseness of  $D(L)$  in  $H$  would then imply

$$g_1 = g_2$$

Since such  $g$  are unique we may define

$$L^*v = g \quad \forall v \in D(L^*)$$

Other than knowing  $0 \in L^*$  we don't know the size of the domain  $D(L^*)$

## Self Adjoint and Symmetric Operators

Authors differ on how to define symmetric and self adjoint. Here I present that of Reed and Simon (I Functional Analysis)

Definition:  $L$  densely defined on  $H$  is symmetric if

$$(i) \quad D(L) \subset D(L^*)$$

$$(ii) \quad Lu = L^*u, \quad \forall u \in D(L)$$

Most applied mathematicians say  $L$  is formally self adjoint if (ii) holds on any subset of  $H$  (dense)

Defn:  $L$  densely defined on  $H$  is self adjoint if

$$(i) \quad D(L) = D(L^*)$$

$$(ii) \quad Lu = L^*u \quad \forall u \in D(L)$$

Self adjoint is a big huge deal.

Symmetric operators are usually extendable to self adjoint ones and is technical detail in functional analysis.



## EXAMPLE      Periodic Boundary conditions

$$Lu \equiv u''$$

$$D(L) \equiv \{u \in C^2[0, \pi] : u(0) = u(\pi), u'(0) = u'(\pi)\}$$

Again we integrate by parts

$$\begin{aligned}\langle Lu, v \rangle &= \int_0^\pi u'' v \, dx \\ &= u'v \Big|_0^\pi - \int_0^\pi u'v' \, dx \\ &= u'v \Big|_0^\pi - uv' \Big|_0^\pi + \int_0^\pi u \underbrace{v''}_{L^*v} \, dx\end{aligned}$$

Owing to periodicity of  $u, u'$  at  $x=0, \pi$

$$= u'(0)(v(\pi) - v(0)) - u(0)(v'(\pi) - v'(0)) + \langle u, L^*v \rangle$$

must vanish since  $u'(0)$   
and  $u(0)$  are arb.

yields to the adjoint operator

$$L^*v = v''$$

$$D(L^*) = \{v \in C^2[a, b] : v(0) = v(\pi), v'(0) = v'(\pi)\}$$

Since  $Lu = L^*u \quad \forall u \in D(L^*) = D(L)$ ,  $L$  is self adjoint

Side note,  $u(x) = \text{const}$  is in  $D(L)$ . But  $u(x) = k \Rightarrow Lu = 0$  hence  $N(L)$  nonempty and  $Lu = f$  can't have unique soln

EXAMPLE  $H = L^2$

$$Lu \equiv u'' + u'$$

$$D(L) \equiv \{u \in C^2[0,1] : u(0) = 0, u'(1) = 0\}$$

To find  $L^*$  we must find an  $L^*$  such that

$$\langle Lu, v \rangle = \langle u, L^*v \rangle$$

for an appropriate domain  $D(L^*)$

$$\langle Lu, v \rangle = \int_0^1 (u'' + u')v \, dx$$

$$\langle Lu, v \rangle = \underbrace{(u' + u)v \Big|_0^1 - uv' \Big|_0^1}_{\text{need this to vanish } \forall u \in D(L)} + \int_0^1 (v'' - v')u \, dx$$

$L^*v$

Since  $u(0) = 0, u'(1) = 0$ , algebra yields

$$\langle Lu, v \rangle = u'(0)v(0) + \underbrace{(v(1) - v'(1)) \cdot u(1)}_{\text{need this to vanish}} + \int_0^1 (v'' - v')u \, dx$$

Domain  $D(L^*)$  must have indicated terms vanish.  
Thus

$$\langle Lu, v \rangle = \langle u, L^*v \rangle$$

for

$$L^*v \equiv v'' - v'$$

$$D(L^*) = \{u \in C^2[0,1] : v(0) = 0, v(1) - v'(1) = 0\}$$

Note  $L^* \neq L$  and  $D(L^*) \neq D(L)$

EXAMPLE Integral operator (Hilbert Schmidt)

$$Lu \equiv \int_0^1 \sin(x+2y) u(y) dy$$

is defined on  $D(L) = H = L^2[0, 1]$ . Since kernel  $k(x, y) \equiv \sin(x+2y)$  is smooth on compact domain

$$\begin{aligned} \langle Lu, v \rangle &= \int_0^1 v(x) \int_0^1 \sin(x+2y) u(y) dy dx \\ &= \int_0^1 \int_0^1 \sin(x+2y) v(x) u(y) dy dx \\ &= \int_0^1 u(y) \underbrace{\int_0^1 \sin(x+2y) v(x) dx}_{L^*v} dy \\ &= \langle u, L^*v \rangle \end{aligned}$$

if

$$L^*v \equiv \int_0^1 \sin(x+2y) v(x) dx$$

Note  $D(L^*) = H$  and

$$L^*u \neq Lu$$

not self adjoint.

EXAMPLE Laplacian on  $H = L^2(\Omega)$ ,  $\Omega \subset \mathbb{R}^3$  compact

$$Lu \equiv \nabla^2 u$$

$$D(L) \equiv \{u \in C^2(\Omega) : u|_{\partial\Omega} = 0\}$$

By applying Divergence Theorem to  $\vec{F} = v \nabla u - u \nabla v$  we get Green's identity (second):

$$\int_{\Omega} (v \nabla^2 u - u \nabla^2 v) dx = \int_{\partial\Omega} (v \frac{\partial u}{\partial n} - u \frac{\partial v}{\partial n}) ds$$

Written another way

$$\langle Lu, v \rangle = \int_{\partial\Omega} v \frac{\partial u}{\partial n} ds - \int_{\partial\Omega} u \frac{\partial v}{\partial n} ds + \int_{\Omega} u \nabla^2 v dx$$

$\uparrow$   
must vanish  
 $\forall u \in D(L)$

$\nearrow$   
 $u \in D(L)$

$\int_{\Omega} u \nabla^2 v dx = \int_{\Omega} u (L^* v) dx$

Conclude  $\langle Lu, v \rangle = \langle u, Lv \rangle \quad \forall u \in D(L), v \in D(L^*)$   
if

$$L^* v = \nabla^2 v$$

$$D(L^*) = D(L)$$

hence  $L$  is self adjoint

EXAMPLE

Volterra integral operator

$$Lu \equiv \int_0^x u(t) dt$$

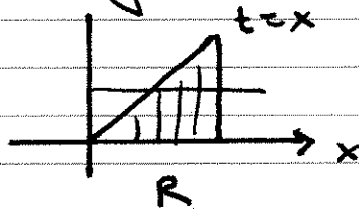
$$D(L) = L^2[0, 1]$$

Then

$$\langle Lu, v \rangle = \int_0^1 \left( \int_0^x u(t) dt \right) v(x) dx$$

$$= \int_0^1 \int_0^x u(t) v(x) dt dx = \iint_R u(t) v(x) dx dt$$

Integration region:



$$t < x < 1$$

$$0 < t < 1$$

Thus

$$\langle Lu, v \rangle = \int_0^1 u(t) \left( \int_t^1 v(x) dx \right) dt$$

and

$$L^*v = \int_x^1 v(t) dt$$

Clearly  $L^* \neq L$

EXAMPLE Let  $H = L^2[0,1]$  with its standard inner product  $\langle \cdot, \cdot \rangle$ . If we let  $H = H_1 \oplus H_1$ , the inner product for this Hilbert space is

$$\langle\langle u, v \rangle\rangle = \langle u_1, v_1 \rangle + \langle u_2, v_2 \rangle$$

where  $u = (u_1, u_2)$  etc.

Consider  $L: H \rightarrow H$  defined by

$$Lu = \frac{d^2 u}{dx^2} + Au \quad A \in \mathbb{R}^{2 \times 2}$$

$$D(L) = \{u \in H : Lu \text{ defined, } u_k(0) = u_k(1) = 0\}$$

Find adjoint

$$\langle Lu, v \rangle = \langle u_1'' + a_{11}u_1 + a_{12}u_2, v_1 \rangle + \langle u_2'' + a_{21}u_1 + a_{22}u_2, v_2 \rangle$$

With some foresight  $D(L^*) = D(L)$  so integrating by parts

$$\langle Lu, v \rangle = \langle u_1, v_1'' \rangle + \langle \underline{a_{11}u_1 + a_{12}u_2}, v_1 \rangle + \langle u_2, v_2'' \rangle + \langle \underline{a_{21}u_1 + a_{22}u_2}, v_2 \rangle$$

underlined part from  $\langle Au, v \rangle$ . Claim this equals  $\langle u, A^T v \rangle$  so that

$$\langle Lu, v \rangle = \langle \frac{d^2 v}{dx^2} + A^T v, u \rangle$$

implies

$$L^* v = \frac{d^2 v}{dx^2} + A^T v$$

## Operators: definitions and properties

In the following definitions and theorems,  $H$  is a Hilbert space though many definitions hold for normed linear spaces.

**Sets:** Let  $S \subset H$  be some set. Then

(C1) $S$ bounded	$\Leftrightarrow \exists M > 0$ s.t. $\ x\  \leq M, \forall x \in S$
(C2) $S$ compact	$\Leftrightarrow$ Every sequence $\{x_n\} \subset S$ contains a convergent subsequence $\{x_{n_k}\}$ which converges to $x \in S$
(C3) $S$ bounded	$\not\Leftrightarrow S$ compact
(C4) $S$ (sequentially) compact	$\Rightarrow S$ closed and bounded
(C5) $S \equiv \{x \in H : \ x\  \leq 1\}$ compact	$\Rightarrow \dim(H) < \infty$

**Definition: Bounded Operator** An operator  $L : H \rightarrow H$  is bounded if there exists some  $M > 0$  such that

$$\|Lx\| \leq M \|x\| \quad \forall x \in H \tag{1}$$

If  $L$  is not bounded then  $L$  is unbounded.

**Definition: Compact Operator** An operator  $L : H \rightarrow H$  is compact if it maps bounded sets into compact sets. Equivalently, for every bounded  $\{x_n\} \subset H$  the sequence  $\{Lx_n\}$  has a convergent subsequence.

**Definition: Adjoint Operator** Let  $L : D(L) \subset H \rightarrow H$ . The operator  $L^*$  defined on  $D(L^*)$  is the adjoint of  $L$  if

$$\langle Lu, v \rangle = \langle u, L^*v \rangle \quad , \quad \forall u \in D(L), \forall v \in D(L^*)$$

Further we say  $L$  is self adjoint if

$$\begin{aligned} Lu &= L^*u \quad , \quad \forall u \in D(L) \\ D(L) &= D(L^*) \end{aligned}$$

$L$  is formally self adjoint if  $Lu = L^*u$  on some set dense subset of  $H$ .

The Reisz representaion Theorem assures that bounded operators always have adjoints. Since compact operators are bounded the same applies to them.

### Theorems for Compact and Bounded linear operators

- (1)  $L$  bounded  $\Leftrightarrow L$  continuous
- (2)  $L$  bounded  $\Leftrightarrow L(\lim_{n \rightarrow \infty} u_n) = \lim_{n \rightarrow \infty} Lu_n$
- (3)  $L$  bounded and  $R(L)$  closed  $\Leftrightarrow R(L) \perp N(L^*)$
- (4)  $L$  compact  $\Rightarrow L$  bounded
- (5)  $L$  linear,  $\dim(R(K)) < \infty \Rightarrow L$  compact
- (6)  $L$  bounded,  $\{\phi_n\}_{n=1}^{\infty}$  orthonormal  $\Rightarrow \lim_{N \rightarrow \infty} L\phi_n = 0$
- (7)  $L_1, L_2$  bounded  $\Rightarrow L_1 + L_2$  bounded
- (8)  $L_1, L_2$  compact  $\Rightarrow L_1 + L_2$  compact
- (9)  $L_n$  compact,  $\|L_n - L\|_{op} \rightarrow 0 \Rightarrow L$  compact

**Definition: Operator norm** For any bounded operator  $L : H \rightarrow H$  we define the operator norm as

$$\|L\|_{op} \equiv \max_{\|x\|=1} \|Lx\| = \max_{x \neq 0} \frac{\|Lx\|}{\|x\|} \quad (2)$$

Theorems above imply that the space of bounded (and compact) operators are normed linear spaces. Completeness can also be shown.



### Self Adjoint Compact Operators

The classic example of a self adjoint compact operator is the Hilbert-Schmidt operator  $L$  defined on  $L^2(\Omega)$  where  $\Omega \subset \mathbb{R}^n$  is compact.

$$Lu = \int_{\Omega} k(x, y)u(y) dy$$

This is well defined if kernel  $k$  is square integrable on  $\Omega^2$  and is self adjoint if

$$k(x, y) = k(y, x)$$

Below  $L$  is a self adjoint compact operator on  $H$  with eigenvalue-vector pairs:

$$L\phi_n = \lambda_n\phi_n \quad , \quad \phi_n \neq 0$$

- a) All the eigenvalues of  $L$  are real
- b) If  $\lambda \neq 0$  then the associated eigenspace has  $\dim E_{\lambda} < \infty$
- c) If  $\lambda_1 \neq \lambda_2$  then  $\langle u_1, u_2 \rangle = 0$ .
- d) Either  $L$  has a finite number of eigenvalues or a countable number with  $\lambda_n \rightarrow 0$ .
- e)  $\exists$  a sequence  $\{\phi_n\}$  of orthonormal eigenfunctions of  $L$  that are a basis for  $H$ <sup>1</sup>

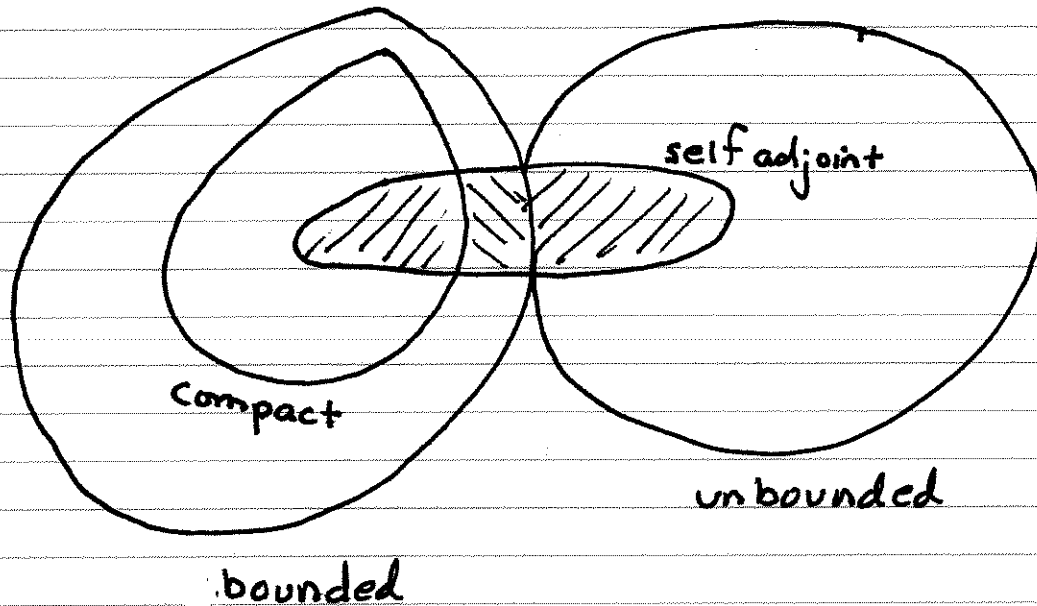
Note that  $\lambda = 0$  is very possible hence  $N(L)$  may be nonempty. In the case of “degenerate” integral operators one can even have  $\dim N(L) = \infty$ . Statement c) above is not true for finite dimensional Hilbert spaces. It applies if  $\dim(H) = \infty$  regardless of whether  $N(L)$  is empty or not. The last fact e) is very important. In practice one uses Gram Schmidt to orthogonalize bases for each eigenspace  $E_{\lambda}(L)$ . Then, if  $L$  is a self adjoint compact operator then the eigenvalues form an orthonormal basis for  $H$  and one can always write:

$$f = \sum_{n=1}^{\infty} \langle f, \phi_n \rangle \phi_n \quad , \quad \forall f \in H$$

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<sup>1</sup>Reed and Simon, Functional Analysis I, pg 203

## Operator Space nestings



A separate issue of importance in Functional analysis is whether:

$R(L)$  closed

$R(L)$  not closed.

## Fredholm Alternative (Bounded $L$ , closed $R(L)$ )

Let  $L: H \rightarrow H$  be bounded and  $R(L)$  closed.

$$Lu = f \text{ has soln } \Leftrightarrow \langle f, v \rangle = 0 \quad \forall v \in N(L^*)$$

Proof. First let  $Lu = f$ . Then

$$\langle Lu, v \rangle = \langle u, L^*v \rangle = \langle f, v \rangle = 0 \quad \forall v \in N(L^*)$$

For the opposite implication we note that since  $R(L)$  is closed "closed"

$$H = R(L) \oplus R(L)^\perp$$

Then  $f = f_r + f_r^\perp$  where  $f_r \in R(L)$ .

$$\begin{aligned} \text{"bounded"} \left( \begin{aligned} \langle f_r^\perp, Lz \rangle &= 0 & \forall z \in H \\ \langle L^*f_r^\perp, z \rangle &= 0 & \forall z \in H \end{aligned} \right. \end{aligned}$$

so that  $f_r^\perp \in N(L^*)$

$$\langle v, f \rangle = 0 \quad \forall v \in N(L^*)$$

$$\langle f_r^\perp, f_r^\perp + f_r \rangle = 0$$

$$\|f_r^\perp\|^2 = 0$$

$$f_r^\perp = 0$$

Conclude  $f = f_r + \cancel{f_r^\perp} = f_r$  hence  $f \in R(L)$  /

## EXAMPLE Integral equation

$$(1) \quad u(x) = f(x) + \lambda \int_0^1 xy u(y) dy$$

Define the (compact, bounded) operator

$$Ku \equiv \int_0^1 xy u(y) dy$$

the (1) is

$$(2) \quad (I - \lambda K)u = f$$

The operator on the left is bounded (and closed range... technical) so solvability of (1) depends on (F.A.Th.)  $N(I - \lambda K)^*$   
Clearly

$$(I - \lambda K)^* = (I - \lambda K)$$

Suffice to find basis for  $N(I - \lambda I) \Leftrightarrow$

$$u(x) = \lambda x \int_0^1 u(y) dy = ax \quad a \in \mathbb{R}$$

i.e.  $N(I - \lambda I) = \text{span}\{x\}$ .

Fred. Alt. Theorem implies (1) has a soln  $\Leftrightarrow$

$$\int_0^1 x f(x) dx = 0$$

## EXAMPLE Integral Eqn perturbed eigenvalues

$$Ku \equiv (K_0 + \varepsilon K_1)u$$

$$Ku \equiv \int_0^1 xyu(y)dy + \varepsilon \int_0^1 u(y)dy$$

### Unperturbed eigenvalue problem

$$(K_0 - \lambda_0 I)u_0 = 0$$

where

$$\lambda_0 u_0 = \int_0^1 xyu(y)dy = ax \quad a \in \mathbb{R}$$

Hence  $u_0(x)$  proportional to  $x$ . wlog  $u_0(x) = x$

$$\lambda_0 = \int_0^1 y^2 dy = \frac{1}{3}$$

Conclude  $\lambda_0 = \frac{1}{3}$  is the (nonzero) e-value

$$\lambda_0 = \frac{1}{3} \quad N(K_0 - \lambda_0 I) = \text{span}\{x\}$$

### Perturbed eigenvalue problem.

Assume the expansions

$$\lambda = \lambda_0 + \varepsilon \lambda_1 + \dots$$

$$u = u_0 + \varepsilon u_1 + \dots$$

and use F.A.Th. to find  $\lambda_1$ .

Collect like powers of  $\varepsilon$  in

$$(\mathbb{K}_0 + \varepsilon \mathbb{K}_1)(u_0 + \varepsilon u_1 + \dots) = (\lambda_0 + \varepsilon \lambda_1 + \dots)(u_0 + \varepsilon u_1 + \dots)$$

yields

$$(\mathbb{K}_0 - \lambda_0 \mathbb{I}) u_0 = 0$$

$$(1) \quad (\mathbb{K}_0 - \lambda_0 \mathbb{I}) u_1 = \lambda_1 u_0 - \mathbb{K}_1 u_0 \equiv f$$

Apply F.A.T. to (1). Note  $(\mathbb{K}_0 - \lambda_0 \mathbb{I})^* = (\mathbb{K}_0 - \lambda_0 \mathbb{I})$   
so

$$\langle \lambda_1 u_0 - \mathbb{K}_1 u_0, v \rangle = 0 \quad \forall v \in N(\mathbb{K}_0 - \lambda_0 \mathbb{I})$$

Previously found  $N(\mathbb{K}_0 - \lambda_0 \mathbb{I}) = \text{span}\{x\}$

$$\langle \lambda_1 x - \mathbb{K}_1(x), x \rangle = 0$$

$$\langle \lambda_1 x - \frac{1}{2}, x \rangle = 0$$

$$\lambda_1 \|x\|^2 - \frac{1}{2} \langle 1, x \rangle = 0$$

$$\frac{1}{3} \lambda_1 - \frac{1}{2} \cdot \frac{1}{2} = 0$$

$$\boxed{\lambda_1 = \frac{3}{4}}$$

## Fredholm Alternative ( $D(L)$ dense in $H$ )

$$Lu = f \text{ has soln } \Rightarrow \langle f, v \rangle = 0 \quad \forall v \in N(L^*)$$

Pf: Trivial. Denseness implies  $L^*$  is uniquely defined on some  $D(L^*)$  though it does not state how "big"  $D(L^*)$  is. Then

$$\langle Lu, v \rangle = \langle u, L^*v \rangle = \langle f, v \rangle \quad \square$$

Indicated terms vanish  $\forall v \in D(L^*)$

Remark: In a compact form

$$f \in R(L) \Rightarrow f \in N(L^*)^\perp$$

Non equivalence does not hinder the usefulness as "necessary" conditions for solns may still be derived.

EX Helmholtz on rectangle  $\Omega = [0, \pi]^2$

$$Lu \equiv \nabla^2 u + 5u$$

$$D(L) = \{u \in C^2(\Omega) : u|_{\partial\Omega} = 0\}$$

$L$  is "a" Helmholtz operator. Readily verified with Green's identity that  $L$  is self adjoint.

$$(1) \quad \nabla^2 u + 5u = f \quad u \in D(L)$$

where  $f = \sin(2x)$ . Question is, does (1) have a solution? Has a soln only if  $f \in N(L^*)^\perp$ .

One can verify  $\mathcal{N} \equiv \text{span}\{\sin 2x \overset{v_1}{\sin y}, \sin x \overset{v_2}{\sin 2y}\} \subset N(L^*)$   
Hard to prove but  $\mathcal{N} = N(L) = N(L^*)$ .

Fred. Alt Theorem implies  $f$  must be orthogonal to both  $v_1$  and  $v_2$ .  
Calculations show

$$\langle v_1, f \rangle = \pi \neq 0$$

$$\langle v_2, f \rangle = 0$$

hence (1) has no solution if  $f = \sin(2x)$ .  
Recall too in the above

$$\langle u, v \rangle = \int_0^\pi \int_0^\pi u(x, y) v(x, y) dy dx$$



EX Helmholtz on rectangle  $\Omega \equiv [0, \pi]^2$

$$Lu \equiv \nabla^2 u + u$$

where the domain for  $L$  is

$$D(L) = \{u \in C^2(\Omega) : u|_{\partial\Omega} = 0\}$$

$L$  is a Helmholtz operator. It's not hard to show that  $L$  is self adjoint.

Consider the problem

$$(1) \quad \nabla^2 u + u = f$$

not in  $D(L)$   
 $u|_{\partial\Omega} = 0$

where  $f(x,y) = \cos(2x)$ .

Eqn (1) has a solution only if  $f \in N(L^*)^\perp$  by F.A.Thm.

One can verify  $T \equiv \text{span}\{\sin x, \sin y\} \subset N(L^*)$ .  
Actual equality by hand to prove.

$$\left. \begin{aligned} \langle \sin x, f \rangle &= -\frac{2\pi}{3} \\ \langle \sin y, f \rangle &= 0 \end{aligned} \right\} \text{not both zero}$$

Hence (1) does not have a soln if  $f = \cos(2x)$

$$\langle u, v \rangle \equiv \int_0^\pi \int_0^\pi u(x,y)v(x,y) dy dx$$

## Resolvents $(I - \lambda K)^{-1}$

Let  $K$  be compact and self adjoint.  
Seek a series solution to

$$(1) \quad (I - \lambda K)u = f$$

Such a solution is unique  $\Leftrightarrow N(I - \lambda K) = 0$   
in which case one might write

$$(2) \quad u = (I - \lambda K)^{-1} f$$

The operator on the right is called  
the resolvent operator

An example of such an equation  
would be the following integral  
eqn

$$u(x) - \lambda \int_a^b k(x,y)u(y)dy = f(x)$$

where the Hilbert Schmidt kernel  
is symmetric

$$k(x,y) = k(y,x)$$

This assures the operator

$$Ku \equiv \int_a^b k(x,y)u(y)dy$$

is self adjoint and has a complete  
set of orthonormal eigen functions

$$(2) \quad K\phi_n = \mu_n \phi_n \quad \langle \phi_i, \phi_j \rangle = \delta_{ij}$$

To solve we let

$$(3) \quad u = f + g$$

in (1) and rearrange to find

$$(4) \quad (I - \lambda K)g = \lambda Kf$$

On account of  $\{\phi_n\}$  being complete ( $K$  comp, s. adj)

$$f = \sum f_n \phi_n \quad f_n = \langle f, \phi_n \rangle$$

$$g = \sum g_n \phi_n \quad g_n = \langle g, \phi_n \rangle$$

So that (4) implies

$$(5) \quad \sum_{n=1}^{\infty} (1 - \lambda \mu_n) g_n \phi_n = \sum_{n=1}^{\infty} \lambda \mu_n f_n \phi_n$$

orthogonality of  $\phi_n \Rightarrow$

$$(6) \quad (1 - \lambda \mu_n) g_n = \lambda \mu_n f_n \quad \forall n$$

Considering  $u = f + g$  and (6) yields  $g_n$

$$(7) \quad u(x) = f(x) + \lambda \sum_{n \geq 1} \frac{\mu_n}{1 - \lambda \mu_n} \langle f, \phi_n \rangle \phi_n = (I - \lambda K)^{-1} f$$

$$u(x) = f(x) + \lambda \int_a^b r(x, y, \lambda) f(y) dy$$

The last written expression only if  $K$  is an integral operator. (7) defines resolvent.