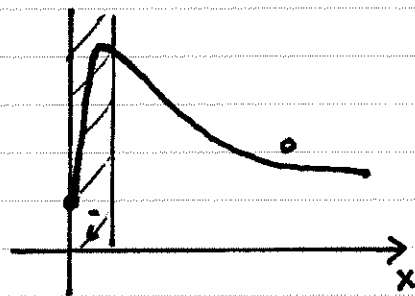


Overview of Linear BVP

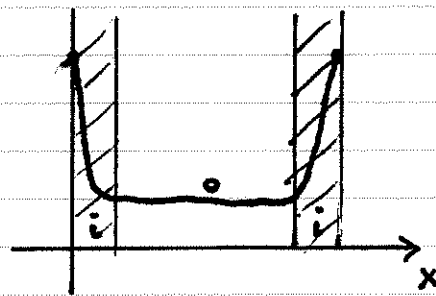
$$(1) \quad \varepsilon y'' + a(x, \varepsilon) y' + b(x, \varepsilon) y = f(x, \varepsilon)$$

$$(2) \quad y(0, \varepsilon) = A \quad y(1, \varepsilon) = B$$

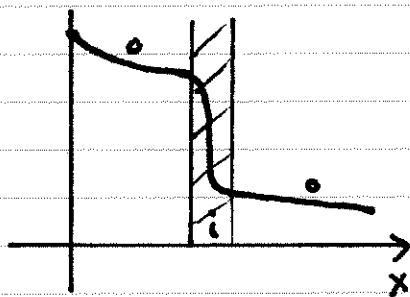
where $\varepsilon \ll 1$. In general $y(x, \varepsilon)$ exhibits singular behavior and has layers



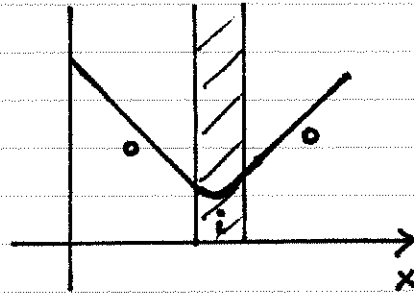
single layer



two layers



interior layer



corner layer

In the layers $\varepsilon y''$ is not small.

δ inner layer

\circ outer layer

Existence of unique solutions

Letting $D = \{y \in C^2[0,1] : y(0) = A, y(1) = B\}$
the problem (1)-(2) may compactly be written

$$(3) \quad L(y) = f \quad y \in D$$

A necessary condition that (3) have a unique soln is that the only solution to

$$(4) \quad L(y) = 0 \quad y(0) = y(1) = 0$$

is $y(x) \equiv 0$. To simplify (4) we use an integrating factor and rewrite

$$y(x) = \exp\left(-\frac{1}{2\varepsilon} \int_0^x a(s) ds\right) w(x)$$

Then $w(x)$ is a soln to

$$(5) \quad \varepsilon w'' - p(x)w = 0$$

$$(6) \quad w(0) = w(1) = 0$$

where

$$(7) \quad p(x) = \frac{a^2}{4\varepsilon} + \frac{a'}{2} - b$$

Then $y(x) \equiv 0$ is the only solution of (4)
iff $w(x) \equiv 0$ is the only soln of (5)-(6)

From (5) we note

$$\varepsilon w w'' - p w^2 = 0$$

Hence

$$\varepsilon \int_0^1 w w'' dx = \int_0^1 p w^2 dx$$

Integrating by parts and using $w(0) = w(1) = 0$ one gets

$$(8) \quad -\varepsilon \int_0^1 (w')^2 dx = \int_0^1 p(x) w^2 dx$$

Notice that if $p(x) < 0$ on some subinterval of $[0, 1]$ this identity could be true for some $w \neq 0$. However, if $p(x) > 0$ on $[0, 1]$ the only soln of (8) is $w = 0$ making the soln of (3) unique. Suff. cond. for uniqueness

$$p(x) = \frac{a^2}{4\varepsilon} + \frac{a'}{2} - b > 0$$

So long as a, a' and b are continuous on $[0, 1]$

$$a(x) \neq 0 \quad \forall x \in [0, 1]$$

assures unique solns as $\varepsilon \rightarrow 0^+$.

Single boundary layer example

$$(1) \quad \varepsilon y'' + a(x)y' + b(x)y = 0 \quad x \in (0, 1)$$

$$(2) \quad y(0) = A \quad y(1) = B$$

$$(3) \quad a(x) > 0 \quad \text{on } [0, 1]$$

We also assume $a, b \in C^1[0, 1]$.

Outer problem

$$y(x, \varepsilon) = y_0(x) + \varepsilon y_1(x) + O(\varepsilon^2)$$

We assume $y_0(x)$ satisfies the right B.C.

$$a(x)y_0' + b(x)y_0 = 0$$

has the solution

$$y_0(x) = B \exp\left(\int_x^1 \frac{b(s)}{a(s)} ds\right)$$

$a \neq 0$

Inner problem

$$y(x, \varepsilon) = Y(\bar{x}, \varepsilon)$$

$$\bar{x} = \frac{x}{\varepsilon}$$

Here there is a boundary layer at $x=0$ of thickness $\delta = O(\varepsilon)$

Differential equation

$$(4) \quad Y'' + a(\epsilon X)Y' + \epsilon b(\epsilon X)Y = 0$$

$$(5) \quad Y(0, \epsilon) = A$$

Using

$$Y(X, \epsilon) = Y_0(X) + \epsilon Y_1(X) + O(\epsilon^2)$$

the leading problem is

$$Y_0'' + a(0)Y_0' = 0, \quad Y_0(0) = A$$

whose solution is

$$(6) \quad Y_0(X) = A + C(1 - e^{-a(0)X})$$

where $C \in \mathbb{R}$ is a constant to be found using matching

Matching

Since $a(0) > 0$ the inner solution Y_0 is bounded as $X \rightarrow \infty$. Consequently the inner and outer solns can be matched:

$$(7) \quad \lim_{x \rightarrow 0^+} y_0(x) = \lim_{X \rightarrow \infty} Y_0(X)$$

Were $a(0) < 0$ the layer must be at $x=1$ since otherwise $Y_0(X) \rightarrow \pm\infty$ can't be matched.

Prandtl matching (7) yields

$$M_0 = \lim_{x \rightarrow 0^+} y_0(x) = B \exp\left(\int_0^1 \frac{b(s)}{a(s)} ds\right) = A + C = \lim_{\bar{x} \rightarrow \infty} \bar{Y}_0(\bar{x})$$

This determines C and the composite soln

$$(8) \quad y_c(x, \varepsilon) \equiv y_0(x) + \bar{Y}_0(\bar{x}) - M_0$$

↓ part that matches in overlap domain.

or

$$y_c(x, \varepsilon) = B \exp\left(\int_x^1 \frac{b(s)}{a(s)} ds\right) + (A - M_0) e^{-a(0)x/\varepsilon}$$

Note $y_c(1, \varepsilon) = B$ whereas $y_c(0, \varepsilon) = A + O(\varepsilon)$.
Also

$$y_c \sim y_0(x) \quad x \text{ fixed}$$

$$y_c \sim \bar{Y}_0(\bar{x}) \quad \bar{x} \text{ fixed}$$

Remarks

$$a(x) > 0 \quad x \in [0, 1] \quad \text{Layer at } x=0$$

$$a(x) < 0 \quad x \in [0, 1] \quad \text{Layer at } x=1$$

$$a(\bar{x}) = 0 \quad \bar{x} \in (0, 1) \quad \text{Interior layer}$$

$$a(x) = 0 \quad x = 0 \quad \text{Corner layer}$$

Depends.

Example Two boundary layers

$$(1) \quad \varepsilon y'' - y = -1 \quad y(0) = 0 \quad y(1) = 2$$

Outer problem

$$y_0(x) \equiv 1$$

can't match either boundary condition hence two boundary layers.

Inner problems

$$y(x, \varepsilon) = \Upsilon(\bar{x}, \varepsilon) \quad \bar{x} = \frac{x - \bar{x}}{\varepsilon^\beta} \quad \bar{x} = 0, 1$$

Substitute into (1) yields

$$\varepsilon^{1-2\beta} \Upsilon'' - \Upsilon = -1$$

Dominant balance $\Rightarrow \beta = \frac{1}{2}$.

Note that $\bar{x} \rightarrow \infty$ for $\bar{x} = 0$ and $\bar{x} \rightarrow -\infty$ for $\bar{x} = 1$ (outer limits)

$$\bar{x} = 1 \quad \Upsilon_0^+(\bar{x}) = c_1^+ e^{\bar{x}} + (1 - c_1^+) e^{-\bar{x}} + 1$$

$$\bar{x} = 0 \quad \Upsilon_0^-(\bar{x}) = -(1 + c_2^-) e^{\bar{x}} + c_2^- e^{-\bar{x}} + 1$$

These satisfy the Bound. Conditions

$$\Upsilon_0^-(0) = 0 \quad \Upsilon_0^+(0) = 2$$

Matching (Prandtl)

The outer solution can match these inner expansions only if:

$$M^+ = \lim_{x \rightarrow 1^-} y_0(x) = \lim_{X \rightarrow -\infty} \Upsilon_0^+(X)$$

$$M^- = \lim_{x \rightarrow 0^+} y_0(x) = \lim_{X \rightarrow \infty} \Upsilon_0^-(X)$$

Considering the expressions for $\Upsilon_0^\pm(X)$ these limits exist, are finite and equal only if

$$c_1^+ = 1 \quad c_2^- = -1$$

Hence

$$\Upsilon_0^+(X) = 1 + e^X \quad X = \frac{x-1}{\varepsilon^{1/2}}$$

$$\Upsilon_0^-(X) = 1 - e^{-X} \quad X = \frac{x}{\varepsilon^{1/2}}$$

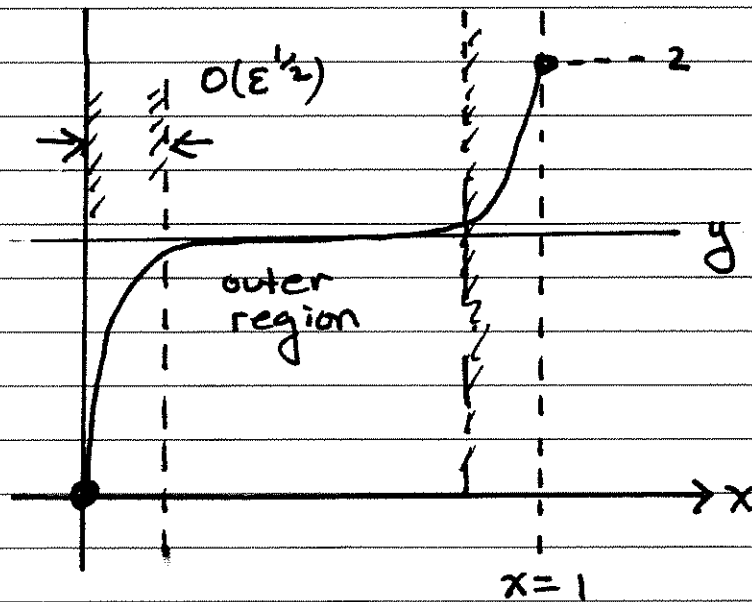
Regardless $M^\pm = 1$, and composite solution

$$y_c(x, \varepsilon) = y_0(x) + \Upsilon_0^+(X) + \Upsilon_0^-(X) - M^+ - M^-$$

Thus

$$y_c(x, \varepsilon) = 1 + \exp\left(\frac{x-1}{\varepsilon^{1/2}}\right) - \exp\left(\frac{x}{\varepsilon^{1/2}}\right)$$

Graph of asymptotic approximation



$$y_c(x, \varepsilon) \sim y_0(x) = 1 \quad \text{in outer region.}$$

EXAMPLE Interior layer

$$\begin{aligned} \epsilon y'' + x y' + 3x^3 y &= 0 & x \in (-1, 1) \\ y(-1) &= 4e & y(1) = 2e^{-1} \end{aligned}$$

Since $a(x) = x$ vanishes at $\bar{x} = 0$ one might expect an interior layer at $\bar{x} \in (-1, 1)$.

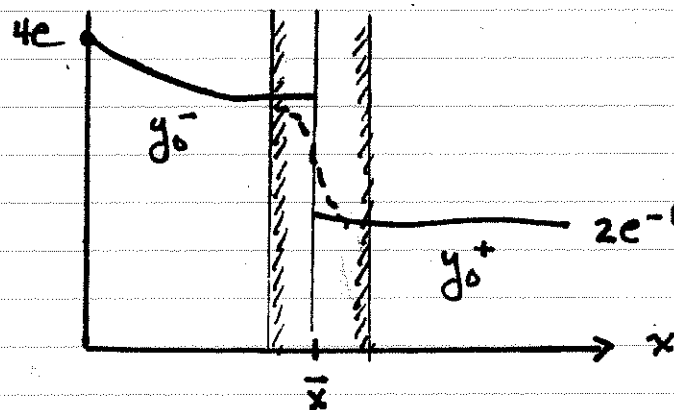
Outer problems

$$x y_0' + 3x^3 y_0 = 0$$

Seek two outer solns in the two outer regions $x > 0$ and $x < 0$.

$$x < 0 \quad y_0^-(x) = 4e^{-x^3} \quad y_0^-(-1) = 4e$$

$$x > 0 \quad y_0^+(x) = 2e^{-x^3} \quad y_0^+(1) = 2e^{-1}$$



Need inner soln ---
to connect
outer solns
in layer

$$\lim_{x \rightarrow \bar{x}} y_0^-(x) = 4 \neq 2 = \lim_{x \rightarrow \bar{x}} y_0^+(x)$$

Inner Problem

$$y(x, \varepsilon) = \Upsilon(\bar{x}, \varepsilon) \quad \bar{x} = \frac{x - \bar{x}}{\delta(\varepsilon)} \quad 0 < \delta \ll 1$$

Here the (interior) layer location is $\bar{x} = 0$ and $\delta(\varepsilon)$ is its thickness to be determined.

$$\Upsilon'' + \frac{\delta^2}{\varepsilon} \bar{x} \Upsilon' + \frac{\delta^5}{\varepsilon} (3\bar{x}^2 \Upsilon) = 0$$

$$\textcircled{1} \sim \textcircled{2} \gg \textcircled{3}$$

Only possible dominant balance (that leads to a Υ_0 which can be matched to both outer solutions)

$$\delta(\varepsilon) = \varepsilon^{1/2}$$

Leading problem

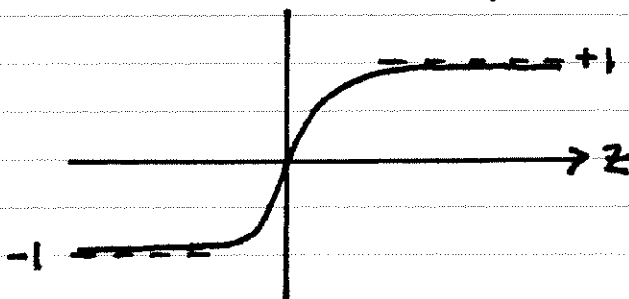
$$\Upsilon_0'' + \bar{x} \Upsilon_0' = 0$$

has the general solution

$$\Upsilon_0(\bar{x}) = C_1 + C_2 \operatorname{erf}\left(\frac{\bar{x}}{\sqrt{2}}\right)$$

where the error function is defined by

$$\operatorname{erf}(z) = \frac{2}{\sqrt{\pi}} \int_0^z e^{-t^2} dt$$

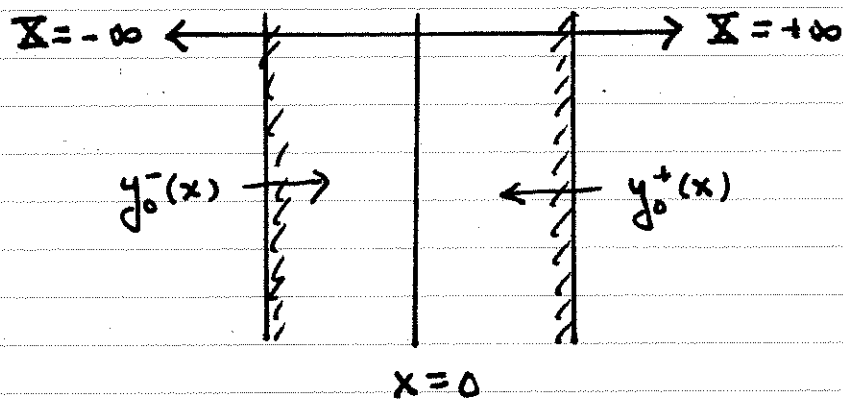


Matching conditions

$$M_+ = \lim_{x \rightarrow 0^+} y_0^+(x) = \lim_{\mathcal{X} \rightarrow \infty} \bar{Y}_0(\mathcal{X})$$

$$M_- = \lim_{x \rightarrow 0^-} y_0^-(x) = \lim_{\mathcal{X} \rightarrow -\infty} \bar{Y}_0(\mathcal{X})$$

Each outer solution must match the (sole) inner solution in their respective overlap domains.



Matching conditions

$$M_+ = 2 = c_1 + c_2$$

$$M_- = 4 = c_1 - c_2$$

$$c_1 = 3, c_2 = -1$$

Hence

$$\bar{Y}_0(\mathcal{X}) = 3 - \operatorname{erf}\left(\frac{\mathcal{X}}{\sqrt{2}}\right)$$

Composite solution (piecewise defined)

$$y_c^\pm(x, \varepsilon) = y_0^\pm(x) + \bar{Y}_0\left(\frac{x}{\sqrt{2\varepsilon}}\right) - M_\pm$$

More explicitly

$$y_c(x, \varepsilon) = \begin{cases} 4e^{-x^3} - \operatorname{erf}\left(\frac{x}{\sqrt{2\varepsilon}}\right) - 1 & x < 0 \\ 2e^{-x^3} - \operatorname{erf}\left(\frac{x}{\sqrt{2\varepsilon}}\right) + 1 & x > 0 \end{cases}$$

For this approximation one can verify

$$\left. \frac{\partial^k y_c}{\partial x^k} \right|_{x=0^-} = \left. \frac{\partial^k y_c}{\partial x^k} \right|_{x=0^+}$$

for $k=0, 1, 2$ but not for $k \geq 3$. Despite the piecewise defn $y_c \in C^2[-1, 1]$

EXAMPLE Varied interior layer thickness

$$\epsilon y'' + x^p y' = 0 \quad y^{(-1)} = 1, y^{(1)} = 3$$

Here $a(x) = x^p$ vanishes at a turning point $\bar{x} = 0$

Outer solutions

Regardless of the region of validity

$$x^p y_0' = 0$$

$$y_0 = c$$

At $x = 0$, $\epsilon y'' = O(x^p y')$ so expect a layer possible.

Inner problem

$$y(x, \epsilon) = \bar{Y}(\bar{x}, \epsilon) \quad \bar{x} = \frac{x - \bar{x}}{\delta(\epsilon)} \quad \bar{x} = 0$$

yields

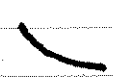


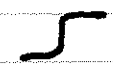
$$\bar{Y}'' + \frac{\delta^{p+1}}{\epsilon} \bar{x}^p \bar{Y}' = 0$$

$$\textcircled{1} \sim \textcircled{2}$$

Balance determines layer thickness

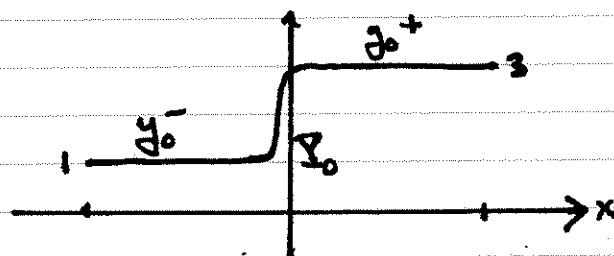
$$\delta(\epsilon) = \epsilon^{\frac{1}{p+1}}$$

Remark: The case $\textcircled{1} \gg \textcircled{2}$ yields $\bar{Y}_0(\bar{x}) = A\bar{x} + B$ which can't be matched to any outer solution.

	p	$\delta(\epsilon)$	$Y_0(x)$	Shape
BL	0	ϵ	$A + B e^{-x}$	
IL	1	$\epsilon^{1/2}$	$A + B \operatorname{erf}\left(\frac{x}{\sqrt{2}}\right)$	
BL	2	$\epsilon^{1/3}$	$A + B \int_0^x \exp(-\frac{1}{3}s^3) ds$	
IL	3	$\epsilon^{1/4}$	$A + B \int_0^x \exp(-\frac{1}{4}s^4) ds$	

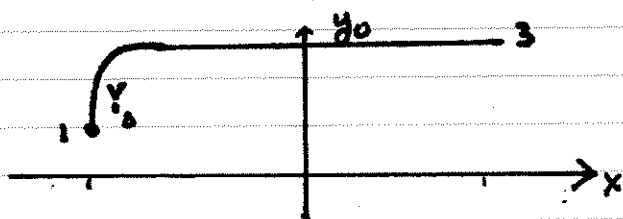
Shapes determine if an interior layer is possible ($p=1, 3$). For $p=0, 2$ inner soln could be a boundary layer.

p odd



interior layer

p even



boundary layer

General Remarks

$$\epsilon y'' + a(x)y' + b(x)y = 0$$

$$a(x) = a'(\bar{x})(x-\bar{x}) + \dots$$

Layer thickness and type depends on T-series of $a(x)$ where $a(\bar{x}) = 0$.

$$a(x) = \sin x \sim x$$

Interior Layer @ $x=0$

$$a(x) = \cos x \sim \frac{(x-\pi/2)^2}{2}$$

B-layer. @ $x=\pi/2$.

EXAMPLE Different B-Layer Thickness

$$(1) \quad \epsilon y'' + x y' - y = x \quad y(0) = y(1) = 1$$

Find an asymptotic approximation having a B-Layer at $x=0$

OUTER SOLUTION

$$x y_0' - y_0 = x$$

has the general solution

$$y_0(x) = x \ln x + C, x$$

Note that $\nexists C$, s.t. $y_0(0) = 1$ satisfies left B.C. Hence a layer must exist at $x=0$ and

$$(2) \quad y_0(x) = x \ln x + x \quad y_0(1) = 1$$

INNER PROBLEM

$$y(x, \epsilon) = Y(X, \epsilon) \quad X = \frac{x}{\epsilon^\beta} \quad \beta > 0$$

yields

$$\underbrace{\epsilon^{1-2\beta} Y''}_{(1)} + \underbrace{X Y'}_{(2)} - Y = \underbrace{\epsilon^\beta X}_{(3)}$$

Seek a dominant balance s.t. $Y_0(X)$ can be matched to $y_0(x)$.

Note ② \Rightarrow ③ for any choice $\beta > 0$.

Were ② \Rightarrow ①, $\mathcal{X} \Upsilon_0' - \Upsilon_0 = 0 \Rightarrow \Upsilon_0(\mathcal{X}) = A \mathcal{X}$
can't satisfy B.C. at $x=0$.

Conclude ① \sim ② and

$$(3) \quad \Upsilon'' + \mathcal{X} \Upsilon' - \Upsilon = \varepsilon^{1/2} \mathcal{X} \quad \beta = \frac{1}{2}$$

Thus $\Upsilon_0(\mathcal{X})$ must solve

$$(4) \quad \Upsilon_0'' + \mathcal{X} \Upsilon_0' - \Upsilon_0 = 0 \quad \Upsilon_0(0) = 1$$

Without symbolic manipulators the solution of (4) can be found by noting

$$\Upsilon_0'''(\mathcal{X}) = \mathcal{X}$$

is a solution and then use variation of parameters to find $\phi(\mathcal{X})$ s.t.

$$\Upsilon_0^{(4)}(\mathcal{X}) = \mathcal{X} \phi(\mathcal{X})$$

is also a solution. After considerable calculations

$$\Upsilon_0(\mathcal{X}) = e^{-\mathcal{X}^2/2} + \mathcal{X} \left(B + \sqrt{\frac{\pi}{2}} \operatorname{erf}\left(\frac{\mathcal{X}}{\sqrt{2}}\right) \right) \quad (4)$$

where constant B found by matching.
Note, in particular, (4) must be bounded as $\mathcal{X} \rightarrow \infty$.

Matching Is there a $B \in \mathbb{R}$ s.t.

$$M = \lim_{x \rightarrow 0} y_0(x) = \lim_{Z \rightarrow \infty} Y_0(Z)$$

Since $\text{erf}(z) = 1 + o(1)$ as $z \rightarrow \infty$, and considering the term (4)

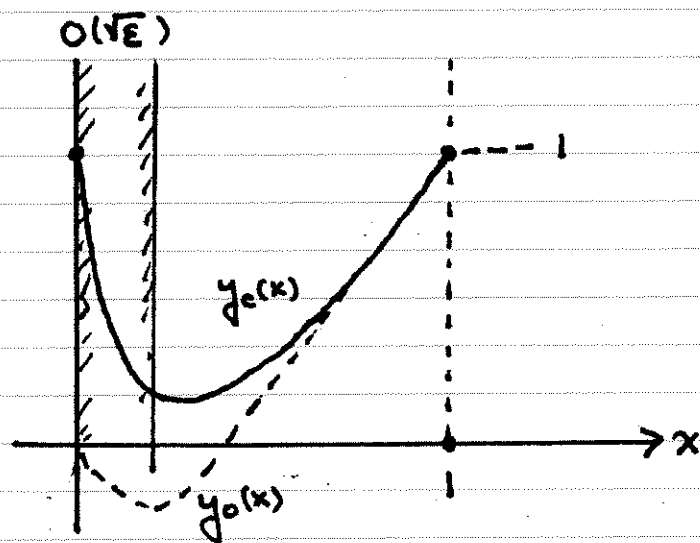
$$B = -\sqrt{\frac{\pi}{2}}$$

so that $M = 0$ and

$$Y_0(Z) = e^{-Z^2/2} + \sqrt{\frac{\pi}{2}} Z \left(\text{erf}\left(\frac{Z}{\sqrt{2}}\right) - 1 \right)$$

Composite solution

$$y_c(x) = y_0(x) + Y_0\left(\frac{x}{\sqrt{\epsilon}}\right)$$



Variation of parameters - addendum

Let $u_1(x)$ be any solution of

$$L(u) = u'' + a(x)u' + b(x)u = 0$$

Seek a second independent solution

$$u_2(x) = \phi(x)u_1(x)$$

Substituting this into $L(u_2) = 0$ gives

$$u_1 \phi'' + (2u_1' + au_1) \phi' + \phi \cancel{L(u_1)} = 0$$

need this to vanish

Thus $v(x) = \phi'(x)$ is a solution of the first order eqn

$$v' + p(x)v = 0$$

$$p(x) = a(x) + \frac{2u_1'}{u_1}$$

which is easily solved so

$$u_2(x) = u_1(x) \int^x v(s) ds$$

Corner Layer Example

$$(1) \quad \varepsilon y'' + xy' - y = 0 \quad x \in (-1, 1)$$

$$(2) \quad y(-1) = 1 \quad y(1) = 2$$

Here $a(x) = x$ and (1)-(2) has a turning point $\bar{x} = 0$

$$a(\bar{x}) = 0$$

$$\bar{x} = 0$$

No boundary layers

$$y(x, \varepsilon) = \bar{Y}(\bar{x}, \varepsilon) \quad \bar{x} \equiv \frac{x - \bar{x}}{\varepsilon}$$

Regardless of $\bar{x} = -1, 1$

$$\bar{Y}_0(\bar{x}) = A + Be^{-\bar{x}\bar{x}}$$

For both $\bar{x} = -1$ and $\bar{x} = +1$ the leading inner solution becomes unbounded in the outer region. For instance, $\bar{x} = -1$

$$\lim_{\bar{x} \rightarrow \infty} \bar{Y}_0(\bar{x}) = \lim_{\bar{x} \rightarrow \infty} (A + Be^{\bar{x}}) = \infty$$

Conclude, no boundary layers.

Outer solution(s)

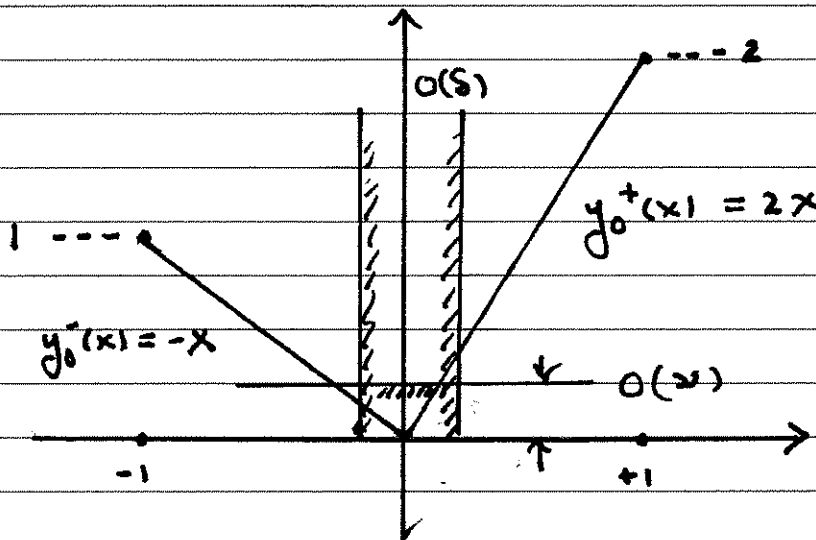
$$x y_0' - y_0 = 0$$

has the general solution $y_0(x) = Ax$.

Independent of the choice of A

$$\lim_{x \rightarrow 0^-} y_0(x) = 0 = \lim_{x \rightarrow 0^+} y_0(x)$$

so we expect a corner layer



$$y_0(x) = \begin{cases} -x & x < 0 \\ 2x & x > 0 \end{cases}$$

A corner layer at $x=0$ exists with

$$x = O(\delta) \quad y = O(\nu)$$

Note: $y_0 \in C(\Omega)$ but $y_0 \notin C^1(\Omega)$ where $\Omega = [-1, 1]$

Inner expansion near corner layer.

Sharp layer near (\bar{x}, \bar{y}) where $\bar{y} = y_0(\bar{x})$

$$\underline{X} = \frac{x - \bar{x}}{\delta(\epsilon)} \quad \delta(\epsilon) \ll 1$$

In the layer the change in y is $O(\nu)$ over $x = O(\delta)$. We don't know either $\delta \ll 1$ or $\nu \ll 1$.

$$(1) \quad y(x, \epsilon) = \Upsilon(\underline{X}, \epsilon) = y_0(\bar{x}) + \nu(\epsilon) \Upsilon_1(\underline{X}) + o(\nu)$$

Here $\Upsilon_0(\underline{X}) = y_0(\bar{x})$ is constant in (1)

The (exact) equation for Υ is

$$(2) \quad \frac{\epsilon}{\delta^2} \Upsilon'' + \underline{X} \Upsilon' - \Upsilon = 0$$

Must choose

$$\delta(\epsilon) = \epsilon^{1/2}$$

Then we get an ODE for $\Upsilon_1(\underline{X})$ which essentially is equivalent to solving the original problem:

$$(3) \quad \Upsilon_1'' + \underline{X} \Upsilon_1' - \Upsilon_1 = 0$$

The general solution of Inner Problem:

$$Y_1(X) = \underbrace{c_1 X}_{\text{unbnd in } X} + c_2 \underbrace{\left(2e^{-\frac{1}{2}X^2} + \sqrt{2\pi} \operatorname{erf}\left(\frac{X}{\sqrt{2}}\right)\right)}_{\text{bounded in } X}$$

Matching to $O(1)$

Prandtl matching condition satisfied

$$\lim_{x \rightarrow 0} y_0(x) = 0 = y_0(\bar{x}) = \lim_{X \rightarrow \pm\infty} Y_0(X)$$

recalling $Y_0(X)$ is constant, i.e. $Y_0(X) \equiv y_0(\bar{x})$

More formally, for $\eta \ll 1$.

$$\lim_{\epsilon \rightarrow 0} \left(y_0(\eta z) - Y_0\left(\frac{\eta z}{\epsilon}\right) \right)$$

$z = \frac{x}{\eta(\epsilon)}$ fixed

$$= \lim_{\epsilon \rightarrow 0} \left(y_0(\eta(\epsilon)z) - y_0(\bar{x}) \right)$$

z fixed

$$= 0$$

Matching to $O(\nu)$ (ν still not known)

outer $y(x, \epsilon) \sim y_0(x) + \epsilon y_1(x) + O(\epsilon^2)$

inner $Y(X, \epsilon) \sim y_0(X) + \nu(\epsilon) Y_1(X) + o(\nu)$

Guess one term outer soln matches two term inner to $O(\nu)$

Intermediate variable $\delta(\epsilon) = \epsilon^{1/2}$

$$z = \frac{x}{\eta(\epsilon)} \quad \delta(\epsilon) \ll \eta(\epsilon) \ll 1$$

Matching conditions are

$$(4) \quad \lim_{\substack{\epsilon \rightarrow 0 \\ z \text{-fixed}}} \frac{(y_0^\pm(\eta z) - \nu(\epsilon) Y_1(\frac{\eta z}{\delta}))}{\nu(\epsilon)} = 0$$

Need a large X approximation of $Y_1(X)$.

From Tables we have

$$(5) \quad \text{erf}(X) = \pm 1 - \frac{1}{\sqrt{\pi} X} e^{-X^2} + \frac{1}{2\sqrt{\pi} X^3} e^{-X^2} + \dots$$

as $X \rightarrow \pm\infty$

Define

$$\frac{M_{01}^+}{v(E)} = \frac{y_0(x) - v(E)Y_1(x)}{v(E)} \quad \underline{\underline{x > 0}}$$

Given previous defns and (5)

$$\frac{M_{01}^+}{v(E)} = \underbrace{\frac{\epsilon^{1/2}}{v} 2x - (\sqrt{2\pi} C_2 + C_1) x}_{\text{need to vanish}} + \underbrace{O(x^{-1} e^{-x^2})}_{\text{T.S.T.}}$$

Conclude

$$\begin{array}{|l} \text{(a)} \quad v(E) = \epsilon^{1/2} \\ \text{(b)} \quad 2 = \sqrt{2\pi} C_2 + C_1 \end{array}$$

A final detail

$$x^2 = \frac{z^2 \eta^2}{\epsilon} = |\ln \epsilon| \Leftrightarrow z\eta = |\epsilon \ln \epsilon|^{1/2}$$

Thus the term above is T.S.T. if

$$|\epsilon \ln \epsilon|^{1/2} \ll \eta \ll 1$$

Also the terms that "matched" above are indicated by \Downarrow . Here, for $x > 0$,

$$M_+ = 2x$$

needed for composite. soln.

Similar expansions for left overlap domain yields

$$(c) \quad -1 = \sqrt{2\pi} C_2 - C_1$$

with the match term

$$M_- = -\mathcal{I}$$

Solving (b)-(c) yields

$$C_1 = \frac{3}{2} \quad C_2 = \frac{\sqrt{2}}{4\sqrt{\pi}}$$

completing inner problem.

Composite Solution

$$\nu(\epsilon) = \epsilon^{1/2}$$

$$y_c(x, \epsilon) = \begin{cases} y_0^+(x) + \nu(\epsilon) (\mathcal{I}_1(\mathcal{X}) - M_+) & x > 0 \\ y_0^-(x) + \nu(\epsilon) (\mathcal{I}_1(\mathcal{X}) - M_-) & x < 0 \end{cases}$$

Moral : Corner layers are bad news from an analysis perspective. Often harder than original problem.

However, outer solns are easy!

EXAMPLE Interior singularity

$$(1) \quad \begin{aligned} \epsilon y'' + 2xy' + (2 + \epsilon x^2)y &= 0 & x \in (-1, 1) \\ y(-1) &= 2 & y(1) = -2 \end{aligned}$$

Solution $y(x)$ is odd

Easy to see the problem (1) including B.C. is invariant under $(x, y) \rightarrow (-x, -y)$.

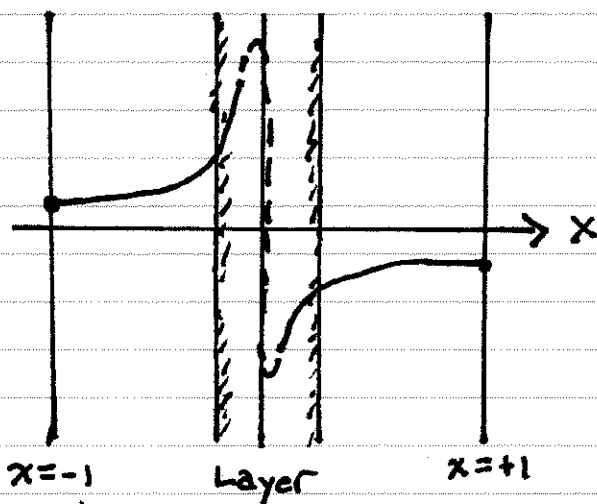
Alternately if $y(x)$ is the solution of (1) for $x > 0$, define $v(z) = -y(-z)$ for $z < 0$. With $z = -x$, can verify $v(z)$ solves (1) for $z < 0$.

Outer Problem $a(x) = 2x$ $\bar{x} = 0$ turning point

$$(2) \quad 2xy_0' + 2y_0 = 0$$

has a general solution $y_0(x) = \frac{A}{x}$, $x \neq 0$.

The outer solution(s) is unbounded in layer



$$y_0^+(x) = -\frac{2}{x} \quad x > 0$$

$$y_0^-(x) = \frac{2}{x} \quad x < 0$$

----- inner soln

Prandtl Matching Fails

Since $y_0(x)$ is unbounded $\Upsilon(\mathcal{X}, \epsilon)$ should be too. Also, Prandtl matching fails since, in particular,

$$\lim_{x \rightarrow 0} y_0(x) \quad \text{D.N.E.}$$

Inner Problem

$$(1) \quad y(x, \epsilon) = \epsilon^{-\alpha} \Upsilon(\mathcal{X}, \epsilon) \quad \mathcal{X} = \frac{x}{\epsilon^\beta}$$

where $\alpha, \beta > 0$ is an assumed asymptotic behavior. That such an inner soln matches for some α, β validates the assumption. For any α

$$(2) \quad \epsilon^{1-2\beta} \Upsilon'' + 2\mathcal{X}\Upsilon' + (2 + o(\epsilon))\Upsilon = 0$$

①

②

③

If ② ~ ③ we repeat the outer problem.
Choose $\beta = \frac{1}{2}$ so ① ~ ② and

$$(3) \quad \Upsilon_0'' + 2\mathcal{X}\Upsilon_0' + 2\Upsilon_0 = 0$$

and

$$y(x, \epsilon) = \epsilon^{-\alpha} \Upsilon_0(\mathcal{X}) + o(\epsilon^{-\alpha})$$

Notice the remainder though asymptotically smaller than $\epsilon^{-\alpha}$, may still be large.

Solution of (3) is

$$Y_0(X) = a e^{-X^2} + b e^{-X^2} \int_0^X e^{t^2} dt$$

Since $y(x, \varepsilon)$ is odd, $Y_0(X)$ must be too, so $a=0$

$$(4) \quad Y_0(X) = b e^{-X^2} \int_0^X e^{t^2} dt \quad b \in \mathbb{R}$$

Matching

wlog need only match $Y_0(X)$ to $y_0^+(x) = -\frac{2}{x}$, $x > 0$
Thus, need large X expansion for $Y_0(X)$.

Use $u = \frac{1}{2t}$ and $v = e^{t^2}$ and integrate (4) by parts.

$$Y_0(X) = \frac{b}{2X} + \underbrace{\frac{1}{2} b \exp(-X^2) \int_0^X \frac{e^{t^2}}{t^2} dt}_{f(X)} + T.S.T$$

Claim $f = O\left(\frac{1}{X^3}\right)$ as $X \rightarrow \infty$. Given claim

$$Y_0(X) = \frac{b}{2X} + O\left(\frac{1}{X^3}\right) \quad \text{as } X \rightarrow \infty$$

Claim is verified by showing

$$\lim_{X \rightarrow \infty} X^3 f(X) = \frac{1}{2}$$

Formal Matching

For $x = \eta z$, $\eta \ll 1$ intermediate variable
need to show

$$(1) \lim_{\substack{\epsilon \rightarrow 0^+ \\ z > 0}} \frac{(y_0(\eta z) - \epsilon^{-\alpha} \mathcal{I}_0(\frac{\eta z}{\epsilon^{1/2}}))}{\nu(\epsilon)} = 0$$

for some $\nu(\epsilon) \ll \epsilon^{-\alpha}$ (in case difference large)

Let M_{00} be numerator of (1). Given $y_0^+(x) = -\frac{2}{x}$

$$M_{00} = \overset{\text{out}}{-\frac{2}{\eta z}} - \epsilon^{-\alpha} \left(\overset{\text{inn.}}{\frac{b}{2x}} + f(x) \right)$$

$$M_{00} = \underbrace{-\frac{2}{\eta z} - \frac{b \epsilon^{1/2-\alpha}}{2\eta z}}_{\text{leading matching terms cancel for all } \nu \text{ if}} - \underbrace{\epsilon^{-\alpha} f\left(\frac{\eta z}{\epsilon^{1/2}}\right)}_{\text{for appropriate restrictions on } \eta(\epsilon) \text{ is small (Difficult)}}$$

$$\alpha = \frac{1}{2} \quad b = -4$$

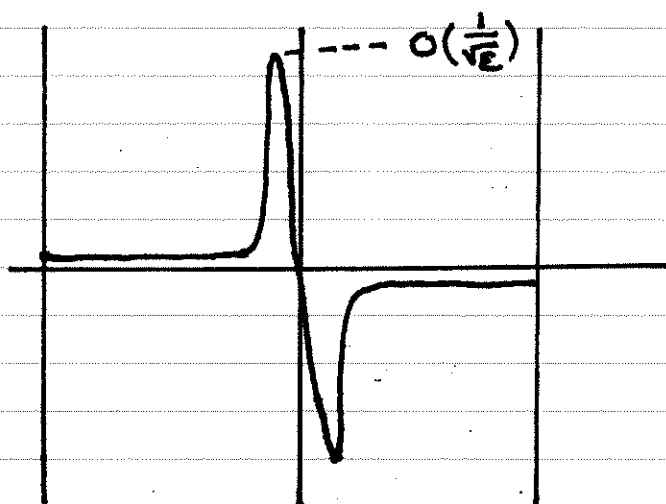
$$\epsilon^{-\alpha} \mathcal{I}_0(x) = -4\epsilon^{-1/2} \exp(-x^2) \int_0^x e^{t^2} dt$$

Composite Solution ($x > 0$ wlog)

$$y_c(x, \epsilon) = y_0(x) + \epsilon^{-1/2} \mathcal{Y}_0(x) - \text{matching terms in overlap region}$$

The matching term is $y_0(x)$ here, so

$$y_c(x, \epsilon) = -4\epsilon^{-1/2} \exp\left(-\frac{x^2}{\epsilon}\right) \int_0^{x/\sqrt{\epsilon}} e^{t^2} dt.$$



dipole pulse.

looks like $\delta'(x)$, delta fn.

Nonlinear Interior Layer example

$$\varepsilon y'' + y(y' + 3) = 0$$

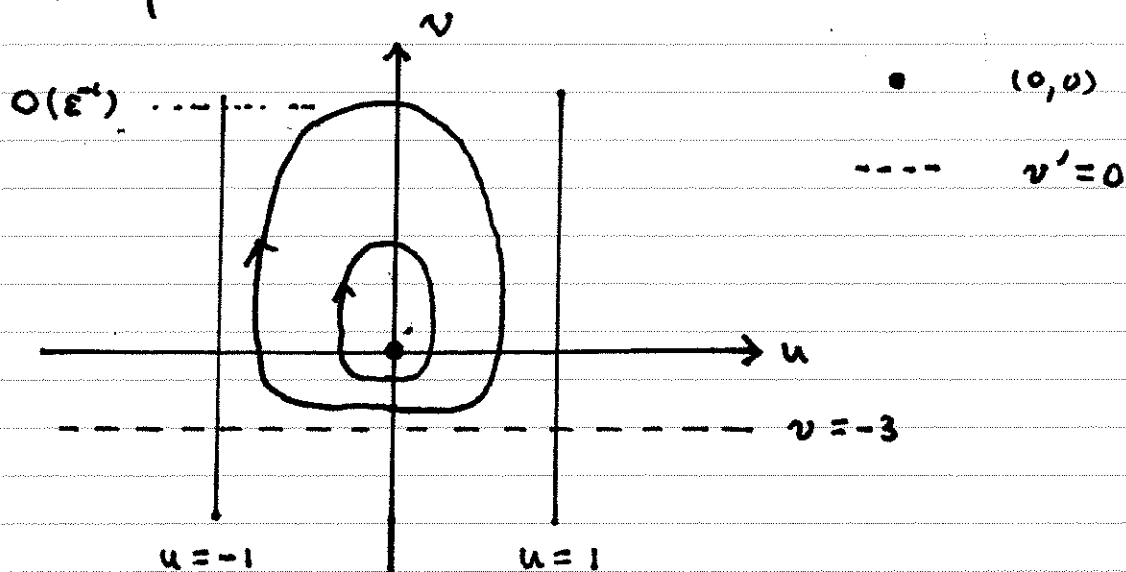
$$y(0) = -1 \quad y(1) = 1$$

As a first order system

$$u' = v$$

$$v' = -\frac{1}{\varepsilon} u(v+3)$$

Has a sole equilibria $(0,0)$ which is a center.
Phase portrait



Boundary condition curves $u = \pm 1$ drawn.

Solution starts on $u = -1$ and ends on $u = +1$
over $x \in (0, 1)$

Outer Solution

$$y_0 (y_0' + 3) = 0$$

$y_0(x) = 0$ not possible since $(0,0)$ equilibria:

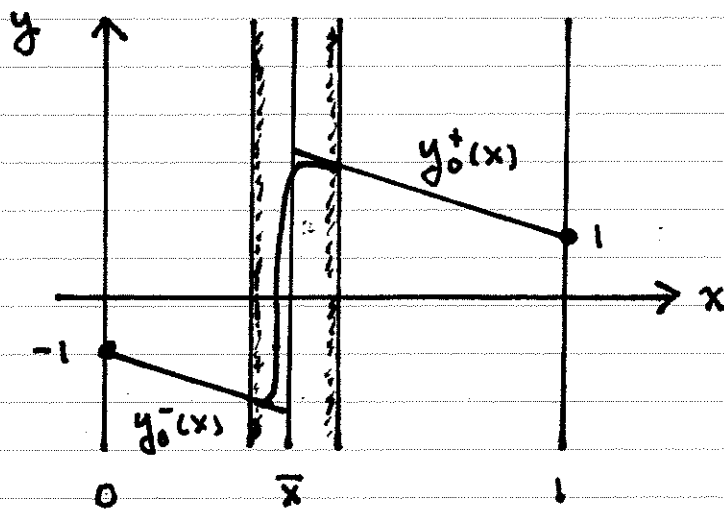
$$y_0^\pm(x) = -3x + C_\pm$$

Using boundary conditions

$$y_0^-(x) = -3x - 1 \quad y_0^-(0) = -1 \quad x < \bar{x}$$

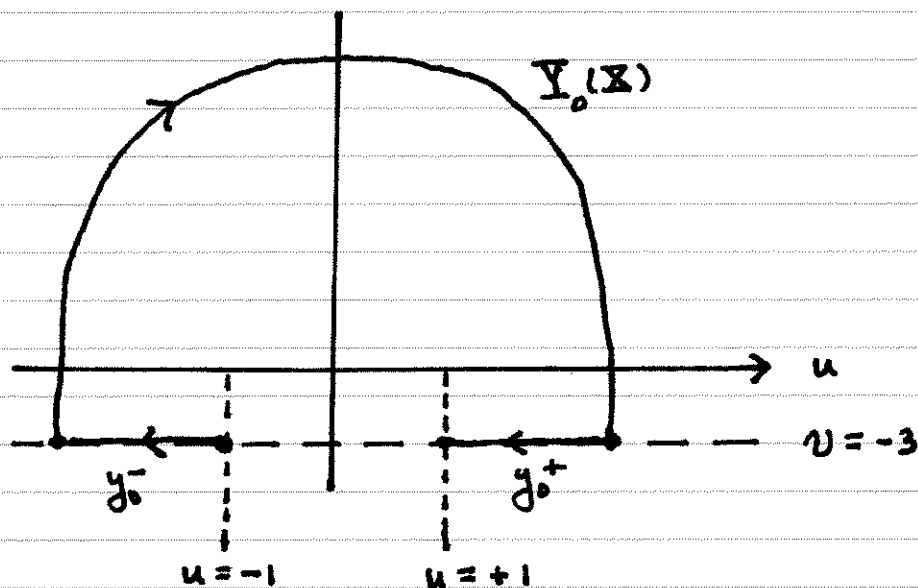
$$y_0^+(x) = -3x + 4 \quad y_0^+(1) = 1 \quad x > \bar{x}$$

Assuming an interior layer at $x = \bar{x}$ where \bar{x} is to be determined.



Inner solution shown with layer thickness $\delta(\epsilon)$

In phase space the solution would look like



Inner Problem

$$y(x, \varepsilon) = Y(X, \varepsilon) \quad X = \frac{x - \bar{x}}{\delta} \quad \delta \ll 1$$

Leads to

$$\frac{\varepsilon}{\delta^2} Y'' + \frac{1}{\delta} Y_0 Y_0' + 3Y = 0$$

$\uparrow \quad \quad \uparrow$
 terms must balance

Thus the B-Layer thickness is $\delta(\varepsilon) = \varepsilon$ and

$$Y_0'' + Y_0 Y_0' = 0$$

To solve inner problem note it is equivalent to

$$\frac{d}{dX} \left(\Upsilon_0' + \frac{1}{2} \Upsilon_0^2 \right) = 0$$

Term in parentheses is constant which leads to a separable first order ODE:

$$\Upsilon_0(X) = c_1 \tanh\left(\frac{c_1}{2}(X - c_2)\right)$$

Matching

$$M_0^- = \lim_{x \rightarrow \bar{x}} y_0^-(x) = \lim_{X \rightarrow -\infty} \Upsilon_0(X)$$

$$M_0^+ = \lim_{x \rightarrow \bar{x}} y_0^+(x) = \lim_{X \rightarrow \infty} \Upsilon_0(X)$$

yields two equations for \bar{x} and c_1 ,

$$M_0^- = -3\bar{x} - 1 = -c_1$$

$$M_0^+ = -3\bar{x} + 4 = c_1$$

Solving

$$\bar{x} = \frac{1}{2} \quad c_1 = \frac{5}{2}$$

Have found layer position but not c_2 . Also

$$M_0^- = -\frac{5}{2} \quad M_0^+ = \frac{5}{2}$$

Value of c_2

$$\Upsilon_0(\mathcal{X}) = \frac{5}{2} \tanh\left(\frac{5}{4}(\mathcal{X} - c_2)\right)$$

$$\Upsilon_0(\mathcal{X}) = \frac{5}{2} \tanh\left(\frac{5}{4}\left(\frac{x - \frac{1}{2}}{\varepsilon} - c_2\right)\right)$$

$$\Upsilon_0(\mathcal{X}) = \frac{5}{2} \tanh\left(\frac{5}{4} \frac{(x - \bar{x}(\varepsilon))}{\varepsilon}\right)$$

where

$$\bar{x}(\varepsilon) = \frac{1}{2} + c_2 \varepsilon$$

so c_2 is $O(\varepsilon)$ correction to layer location.
To leading order we may assume $c_2 = 0$ w.l.o.g.

$$\Upsilon_0(\mathcal{X}) = \frac{5}{2} \tanh\left(\frac{5}{4} \mathcal{X}\right)$$

Composite solution

$$y_c^\pm(x, \varepsilon) = y_0^\pm(x) + \Upsilon_0(\mathcal{X}) - M^\pm \quad (\text{match terms})$$

is piecewise defined

$$y_c(x, \varepsilon) = \begin{cases} -3x - 1 + \frac{5}{2} \tanh\left(\frac{5(x - \frac{1}{2})}{4\varepsilon}\right) + \frac{5}{2} & x < \frac{1}{2} \\ -3x + 4 + \frac{5}{2} \tanh\left(\frac{5(x - \frac{1}{2})}{4\varepsilon}\right) - \frac{5}{2} & x \geq \frac{1}{2} \end{cases}$$