

## 17 (2 March 2009)

### Transport coefficients, cont.

Result: 1<sup>st</sup>-order response to a perturbation  $\mathcal{L}_{\text{per}} = \mathcal{L} + \vec{\lambda} \cdot \vec{\delta}$ , where  $\vec{\delta} = (\delta_1, \dots, \delta_n)$  are commuting derivations and  $\vec{\lambda} = (\lambda_1, \dots, \lambda_n)$  are complex constants.

Then

$$\langle \delta_\nu H \rangle_{\text{per}} = \sum_{\mu} \sigma_{\nu, \mu}^{\delta} \lambda_{\mu} + \mathcal{O}(\lambda^2) \quad (0^{\text{th}}\text{-order response} = 0)$$

defines the first order response (transport) coefficients

$$\begin{aligned} \sigma_{\nu, \mu}^{\delta} &= \text{C-lim}_{t \rightarrow \infty} \tau \left( \rho \int_{\substack{s_0, s_1 \geq 0 \\ s_0 + s_1 = t}} \underbrace{e^{-is_0 \mathcal{L}}}_{\rho} \delta_{\mu} e^{-is_1 \mathcal{L}} \delta_{\nu} H \, ds_0 ds_1 \right) \\ &\quad | \\ &\quad \text{(put this onto } \rho, \text{ but } \rho \text{ is invariant so cross this term off)} \\ &= \text{C-lim}_{t \rightarrow \infty} \tau \left( \int_{\substack{s_0, s_1 \geq 0 \\ s_0 + s_1 = t}} \underbrace{(e^{-is_0 \mathcal{L}}(\rho))}_{\rho} \delta_{\mu} e^{-is_1 \mathcal{L}} \delta_{\nu} H \, ds_0 ds_1 \right) \\ &\quad | \\ &= \text{C-lim}_{t \rightarrow \infty} \tau \left( \delta_{\mu} \int_{\substack{s_0, s_1 \geq 0 \\ s_0 + s_1 = t}} e^{-is_1 \mathcal{L}} \delta_{\nu} H \, ds_0 ds_1 \right) \\ &= 0 - \text{C-lim}_{t \rightarrow \infty} \tau \left( (\delta_{\mu}(\rho)) \int_0^t e^{-is_1 \mathcal{L}} \delta_{\nu} H \, ds_1 \right) \\ &\quad | \\ &\quad \text{because } \tau \text{ is invariant under derivation } (\tau \circ \delta = 0), \text{ then } \tau(\delta_{\mu}(\rho f \text{ etc.})) = 0 \\ &= -\tau(\delta_{\mu}(\rho) \mathcal{D}(\mathcal{L}) \delta_{\nu}(H)) \quad (\text{which may be infinite}) \\ &\quad | \\ &= \lim_{t \rightarrow \infty} \int_0^t e^{-is\omega} \, ds \stackrel{\text{dist'n}}{=} \text{PV} \frac{1}{i\omega} + \frac{\pi}{2} \delta(\omega) =: \mathcal{D}(\omega) \end{aligned}$$

Attention if  $\delta_{\nu}(H)$  is not  $\perp$  to  $\ker \mathcal{L}$ .

Assumption  $T \searrow 0$  and  $E_F \in \text{Gap}(H)$   
Fermi energy

Mathematical assumption is  $\rho = P_F =$  the spectral projection of  $H$  onto states  $\leq E_F$ . This avoids the singularity.

Suppose we have  $\{\psi_i\}_i$  an eigenbasis of  $H$  (in the generalized sense).

So  $H\psi_i = E_i\psi_i$  (that is,  $H|\psi_i\rangle = E_i|\psi_i\rangle$ ).

$$\begin{aligned} \langle \psi_i | \mathcal{L}(\delta_\nu H) | \psi_j \rangle &= \langle \psi_i | (H\delta_\nu H - (\delta_\nu H)H) | \psi_j \rangle \\ &\quad \Big| \\ &\quad \mathcal{L}(\cdot)=[H, \cdot] \\ &= (E_i - E_j)\langle \psi_i | \delta_\nu H | \psi_j \rangle \end{aligned}$$

So,

$$\begin{aligned} \langle \psi_i | \mathcal{L}^n(\delta_\nu H) | \psi_j \rangle &= \langle \psi_i | \mathcal{L}(\mathcal{L}^{n-1}(\delta_\nu H)) | \psi_j \rangle \\ &= (E_i - E_j)\langle \psi_i | \mathcal{L}^{n-1}(\delta_\nu H) | \psi_j \rangle \end{aligned}$$

Hence

$$\langle \psi_i | e^{-is_1\mathcal{L}}\delta_\nu H | \psi_j \rangle = e^{-is_1(E_i - E_j)}\langle \psi_i | \delta_\nu H | \psi_j \rangle$$

So we need that  $E_i \neq E_j$  (to avoid  $\omega = 0$  in the distribution  $\mathcal{D}(\omega)$ ). In that case

$$\langle \psi_i | \lim_{t \rightarrow \infty} \int_0^t e^{-is_1\mathcal{L}} ds_1 \delta_\nu H | \psi_j \rangle = \frac{1}{i(E_i - E_j)} \langle \psi_i | \delta_\nu H | \psi_j \rangle$$

Since  $\rho = P_F$  is a projection and  $\delta_\mu$  is a derivation (using the beautiful formula (\*) from the previous lecture)

$$\delta_\mu P_F = P_F(\delta_\mu P_F)P_F^\perp + P_F^\perp(\delta_\mu P_F)P_F$$

Now

$$\langle \psi_i | \delta_\nu H | \psi_j \rangle \neq 0 \Leftrightarrow ((E_i \leq E_F) \wedge (E_j > E_F)) \vee ((E_j \leq E_F) \wedge (E_i > E_F))$$

Since  $E_F \in (\tilde{E}_0, \tilde{E}_1)$ , a gap in  $\sigma(H)$ , we have

$$|E_i - E_j| \geq |\tilde{E}_1 - \tilde{E}_0|$$

As a consequence,

$$\begin{aligned} & \underset{1=P_F+P_F^\perp}{=} \\ \sigma_{\nu\mu}^\delta &= -\tau(\delta_\mu(P_F)\mathcal{D}(\mathcal{L})\delta_\nu(H)) \\ &= -\tau(P_F\delta_\mu(P_F)(\mathcal{D}\mathcal{L})\delta_\nu(H)) - \tau(P_F^\perp\delta_\mu(P_F)\mathcal{D}(\mathcal{L})\delta_\nu(H)) \\ &= -\tau(P_F\delta_\mu(P_F)(\mathcal{D}(\mathcal{L})\delta_\nu(H))P_F) - \tau(P_F^\perp\delta_\mu(P_F)(\mathcal{D}(\mathcal{L})\delta_\nu(H))P_F^\perp) \\ &= -\tau\left(P_F\delta_\mu(P_F)P_F^\perp\frac{1}{i}\mathcal{L}^{-1}(\delta_\nu H)P_F\right) - \tau\left(P_F^\perp\delta_\mu(P_F)P_F\frac{1}{i}\mathcal{L}^{-1}(\delta_\nu H)P_F^\perp\right) \\ &\stackrel{\text{claim}}{=} -i\tau(P_F\delta_\mu(P_F)\delta_\nu(P_F) - \delta_\nu(P_F)\delta_\mu(P_F)) \\ &\stackrel{\text{or}}{=} -i\tau(P_F\delta_\mu(P_F)P_F^\perp\delta_\nu(P_F)P_F - \delta_\nu(P_F)\delta_\mu(P_F)) \end{aligned}$$

In other words the claim is

$$\begin{aligned} P_F^\perp \mathcal{L}^{-1}(\delta_\nu H) P_F &= -P_F^\perp(\delta_\nu P_F) P_F \quad \text{and} \\ P_F \mathcal{L}^{-1}(\delta_\nu H) P_F^\perp &= P_F(\delta_\nu P_F) P_F^\perp \end{aligned}$$

*Proof.* Hit both sides of the first equation with  $\mathcal{L}$  ( $\mathcal{L}$  leaves  $P_F$  invariant).

$$\begin{aligned} \text{left-hand side: } \mathcal{L}(P_F^\perp \mathcal{L}^{-1}(\delta_\nu H) P_F) &= P_F^\perp(\delta_\nu H) P_F \\ \text{right-hand side: } \mathcal{L}(-P_F^\perp(\delta_\nu P_F) P_F) &= -P_F^\perp[H, \delta_\nu P_F] P_F \\ &\quad \Big| \\ &\quad \mathcal{L}(\cdot) = [H, \cdot] \\ &= -P_F^\perp(\delta_\nu([H, P_F]) - [\delta_\nu H, P_F]) P_F \\ &\quad \Big| \\ &\quad [H, P_F] = 0 \\ &= P_F^\perp((\delta_\nu H) P_F - P_F(\delta_\nu H)) P_F \\ &= P_F^\perp(\delta_\nu H) P_F \quad (\text{QED first equation}) \end{aligned}$$

The proof of the second equation is similar and picks up an extra minus sign.  $\square$

So this is the result

$$\sigma_{\nu\mu}^\delta = i \tau(P_F [(\delta_\nu P_F), (\delta_\mu P_F)])$$

Consequence:  $\sigma_{\nu\mu}^\delta$  is a topological invariant.

*Remark.*  $\sigma_{\nu\mu}^\delta$  is anti-symmetric.

**Example.** QHE in  $\mathbb{R}^2$  (the quantum Hall effect).

$$\begin{aligned} \delta &= (\delta_1, \delta_2) & \delta_\nu &= [\hat{q}_\nu, \cdot] \\ \lambda &= (\lambda_1, \lambda_2) & \lambda_\nu &= e E_\nu \end{aligned}$$

$e$  is the electric charge

$E_\nu$  is the external electric field